Cosmic Rays and Magnetic Fields in Galaxy Formation

Christoph Pfrommer¹

in collaboration with

PhD students: L. Jlassi¹, R. Lemmerz¹, M. Weber¹, J. Whittingham¹
Postdocs: T. Berlok¹, V. Bresci¹, T. Buck², P. Girichidis², L. M. Perrone¹, M. Shalaby¹, M. Sparre^{3,1}, T. Thomas¹, M. Werhahn⁴
Faculty: E. Puchwein¹, R. Pakmor⁴, V. Springel⁴, T. Enßlin⁴

d C

¹AIP Potsdam, ²Heidelberg U, ³U of Potsdam, ⁴MPA Garching *Mike Norman Festschrift*, Kloster Seeon, June 2023



Cosmic ray transport Cosmic rays in galaxy formation Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Puzzles in galaxy formation



Cosmic ray transport Cosmic rays in galaxy formation Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Puzzles in galaxy formation



Christoph Pfrommer Cosmic Rays

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray transport Cosmic rays in galaxy formation Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Puzzles in galaxy formation



Christoph Pfrommer Cosmi

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray transport Cosmic rays in galaxy formation Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Stellar feedback



super wind in M82

NASA/JPL-Caltech/STScI/CXC/UofA

- thermal pressure provided by supernovae or active galactic nuclei?
- radiation pressure and photoionization by massive stars and quasars?
- pressure of cosmic rays (CRs) that are accelerated at supernova shocks?



Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Galactic cosmic ray spectrum



- spans more than 33 decades in flux and 12 decades in energy
- "knee" indicates characteristic maximum energy of galactic accelerators
- CRs beyond the "ankle" have extra-galactic origin



Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Galactic cosmic ray spectrum



- spans more than 33 decades in flux and 12 decades in energy
- "knee" indicates characteristic maximum energy of galactic accelerators
- CRs beyond the "ankle" have extra-galactic origin
- energy density of cosmic rays, magnetic fields, and turbulence in the interstellar gas all similar



Puzzles in galaxy formation Galaxy formation paradigm Cosmic ray population

Review on cosmic ray feedback

The Astronomy and Astrophysics Review manuscript No. (will be inserted by the editor)

Cosmic ray feedback in galaxies and galaxy clusters

A pedagogical introduction and a topical review of the acceleration, transport, observables, and dynamical impact of cosmic rays

Mateusz Ruszkowski^{1,3,*}, Christoph Pfrommer^{2,*}

COSMO



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray transport: an extreme multi-scale problem





Milky Way-like galaxy:

gyro-orbit of GeV CR:

$$r_{
m gal} \sim 10^4 \
m pc$$
 $r_{
m cr} = rac{p_\perp}{e B_{
m uG}} \sim 10^{-6} \
m pc \sim rac{1}{4} \
m AU$

\Rightarrow need to develop a fluid theory for a collisionless, non-Maxwellian component!

Zweibel (2017), Jiang & Oh (2018), Thomas & CP (2019)

Introduction Wave-Cosmic ray transport CR dri Cosmic rays in galaxy formation CR hy

Wave-particle interactions CR driven instabilities CR hydrodynamics

Interactions of CRs and magnetic fields

Cosmic ray



sketch: Jacob & CP



3

Christoph Pfrommer Cosmic Rays and Magnetic Fields in Galaxy Formation

A D > A B >

Wave-particle interactions

Interactions of CRs and magnetic fields



sketch: Jacob & CP



Interactions of CRs and magnetic fields



• electric fields vanish in the Alfvén wave frame: $abla imes {m E} = -rac{1}{c} rac{\partial {m B}}{\partial t}$



Interactions of CRs and magnetic fields



• electric fields vanish in the Alfvén wave frame: $abla imes {m E} = -rac{1}{c} rac{\partial {m B}}{\partial t}$

• work out Lorentz forces on CRs in wave frame: $F_{L} = q \frac{\mathbf{v} \times \mathbf{B}}{C}$



Interactions of CRs and magnetic fields



- electric fields vanish in the Alfvén wave frame: $\nabla \times E = -\frac{1}{c} \frac{\partial B}{\partial t}$
- work out **Lorentz forces on CRs** in wave frame: $F_{L} = q \frac{V \times B}{C}$
- Lorentz force depends on relative phase of CR gyro orbit and wave:
 - sketch: decelerating Lorentz force along CR orbit $ightarrow
 ho_{\parallel}$ decreases
 - phase shift by 180°: accelerating Lorentz force $ightarrow p_{\parallel}$ increases



Interactions of CRs and magnetic fields



sketch: Jacob & CP

 only electric fields can provide work on charged particles and change their energy



Wave-particle interactions

Interactions of CRs and magnetic fields



sketch: Jacob & CP

- only electric fields can provide work on charged particles and change their energy
- in Alfvén wave frame, where E = 0, CR energy is conserved: $p^2 = p_{\parallel}^2 + p_{\perp}^2 = \text{const.}$ so that decreasing p_{\parallel} causes p_{\perp} to increase



Interactions of CRs and magnetic fields



sketch: Jacob & CP

- only electric fields can provide work on charged particles and change their energy
- in Alfvén wave frame, where E = 0, CR energy is conserved: $p^2 = p_{\parallel}^2 + p_{\perp}^2 = \text{const. so that decreasing } p_{\parallel} \text{ causes } p_{\perp} \text{ to increase}$

• this increases the CR pitch angle cosine $\mu = \cos \theta = \frac{B}{|B|} \cdot \frac{p}{|p|}$



Interactions of CRs and magnetic fields



sketch: Jacob & CP

AIP

• CRs resonantly interact with Alfvén waves so that the wavelength equals the gyro-radius:

$$L_{\parallel} = r_{\rm g} = \frac{p_{\perp}c}{qB}$$

Interactions of CRs and magnetic fields



sketch: Jacob & CP

 CRs resonantly interact with Alfvén waves so that the wavelength equals the gyro-radius:

$$L_{\parallel} = r_{g} = rac{p_{\perp}c}{qB}$$

• gyro resonance: $\omega - k_{\parallel} v_{\parallel} = n\Omega = n rac{qB}{\gamma m_i c}$

Doppler-shifted MHD frequency is a multiple n of the CR gyrofrequency

Wave-particle interactions

Interactions of CRs and magnetic fields



CRs resonantly interact with Alfvén waves so that the wavelength equals the gyro-radius:

$$L_{\parallel} = r_{\rm g} = \frac{p_{\perp}c}{qB}$$

$\omega - k_{\parallel} v_{\parallel} = n\Omega = n \frac{qB}{\gamma m_i c}$ gyro resonance:

Doppler-shifted MHD frequency is a multiple *n* of the CR gyrofrequency

AIP

Introduction Wave-part Cosmic ray transport CR driven Cosmic rays in galaxy formation CR hydrod

Wave-particle interactions CR driven instabilities CR hydrodynamics

Cosmic ray streaming and diffusion

- CR streaming instability: Kulsrud & Pearce 1969
 - if v_{cr} > v_a, CR flux excites and amplifies an Alfvén wave field in resonance with the gyroradii of CRs
 - scattering off of this wave field limits the (GeV) CRs' bulk speed ~ v_a
 - wave damping: transfer of CR energy and momentum to the thermal gas





Introduction Wave-partii Cosmic ray transport CR driven i Cosmic rays in galaxy formation CR hydrody

Wave-particle interactions CR driven instabilities CR hydrodynamics

Cosmic ray streaming and diffusion

- CR streaming instability: Kulsrud & Pearce 1969
 - if v_{cr} > v_a, CR flux excites and amplifies an Alfvén wave field in resonance with the gyroradii of CRs
 - scattering off of this wave field limits the (GeV) CRs' bulk speed ~ v_a
 - wave damping: transfer of CR energy and momentum to the thermal gas



 \rightarrow CRs exert pressure on thermal gas via scattering on Alfvén waves



Introduction Wave-partic Cosmic ray transport CR driven in Cosmic rays in galaxy formation CR hydrody

Wave-particle interactions CR driven instabilities CR hydrodynamics

Cosmic ray streaming and diffusion

- CR streaming instability: Kulsrud & Pearce 1969
 - if v_{cr} > v_a, CR flux excites and amplifies an Alfvén wave field in resonance with the gyroradii of CRs
 - scattering off of this wave field limits the (GeV) CRs' bulk speed ~ v_a
 - wave damping: transfer of CR energy and momentum to the thermal gas



 \rightarrow CRs exert pressure on thermal gas via scattering on Alfvén waves

weak wave damping: strong coupling \rightarrow CR stream with waves strong wave damping: less waves to scatter \rightarrow CR diffusion prevails



ヘロト ヘヨト ヘヨト

Introduction Cosmic ray transport

Cosmic rays in galaxy formation

Wave-particle interactions CR driven instabilities CR hydrodynamics

Modes of CR propagation



Wave-particle interactions CR driven instabilities CR hydrodynamics

Modes of CR propagation



Wave-particle interactions CR driven instabilities CR hydrodynamics

Modes of CR propagation



CR vs. radiation hydrodynamics

 capitalize on analogies of CR and radiation hydrodynamics (Jiang & Oh 2018) derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)



> < ≣

< 🗇 🕨

CR vs. radiation hydrodynamics

- capitalize on analogies of CR and radiation hydrodynamics (Jiang & Oh 2018) derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for CR energy and momentum density, $\varepsilon_{\rm cr}$ and $f_{\rm cr}/c^2$

$$\frac{\partial \varepsilon_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{f}_{\rm cr} = -\boldsymbol{w}_{\pm} \cdot \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{v} \cdot \boldsymbol{g}_{\rm Lorentz} + S_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \boldsymbol{f}_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{P}_{\rm cr} = - \qquad \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{g}_{\rm Lorentz} + \boldsymbol{S}_{f}$$

Alfvén wave velocity in lab frame: $w_{\pm} = v \pm v_a$, CR pressure tensor $P_{cr} = P_{cr} \mathbf{1}$, CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

CR vs. radiation hydrodynamics

- capitalize on analogies of CR and radiation hydrodynamics (Jiang & Oh 2018) derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for CR energy and momentum density, $\varepsilon_{\rm cr}$ and $f_{\rm cr}/c^2$

$$\frac{\partial \varepsilon_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{f}_{\rm cr} = -\boldsymbol{w}_{\pm} \cdot \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{v} \cdot \boldsymbol{g}_{\rm Lorentz} + S_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \boldsymbol{f}_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{P}_{\rm cr} = - \qquad \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{g}_{\rm Lorentz} + \boldsymbol{S}_{f}$$

Alfvén wave velocity in lab frame: $w_{\pm} = v \pm v_a$, CR pressure tensor $P_{cr} = P_{cr} \mathbf{1}$, CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

• lab-frame equ's for radiation energy and momentum density, ε and f/c^2 (Mihalas & Mihalas 1984, Lowrie+ 1999):

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot \mathbf{f} = -\sigma_{s} \mathbf{v} \cdot [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_{a}$$
$$\frac{1}{c^{2}} \frac{\partial \mathbf{f}}{\partial t} + \nabla \cdot \mathbf{P} = -\sigma_{s} \quad [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_{a} \mathbf{v}$$



CR vs. radiation hydrodynamics

- capitalize on analogies of CR and radiation hydrodynamics (Jiang & Oh 2018) derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for CR energy and momentum density, $\varepsilon_{\rm cr}$ and $f_{\rm cr}/c^2$

$$\frac{\partial \varepsilon_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{f}_{\rm cr} = -\boldsymbol{w}_{\pm} \cdot \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{v} \cdot \boldsymbol{g}_{\rm Lorentz} + S_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \boldsymbol{f}_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{P}_{\rm cr} = - \qquad \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{g}_{\rm Lorentz} + \boldsymbol{S}_{f}$$

Alfvén wave velocity in lab frame: $w_{\pm} = v \pm v_a$, CR pressure tensor $P_{cr} = P_{cr} \mathbf{1}$, CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

• lab-frame equ's for radiation energy and momentum density, ε and f/c^2 (Mihalas & Mihalas 1984, Lowrie+ 1999):

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot \mathbf{f} = -\sigma_{s} \mathbf{v} \cdot [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_{a}$$
$$\frac{1}{c^{2}} \frac{\partial \mathbf{f}}{\partial t} + \nabla \cdot \mathbf{P} = -\sigma_{s} \quad [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_{a} \mathbf{v}$$

• problem: CR lab-frame equation requires resolving rapid gyrokinetics!

CR vs. radiation hydrodynamics

- capitalize on analogies of CR and radiation hydrodynamics (Jiang & Oh 2018) derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for CR energy and momentum density, $\varepsilon_{\rm cr}$ and $f_{\rm cr}/c^2$

$$\frac{\partial \varepsilon_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{f}_{\rm cr} = -\boldsymbol{w}_{\pm} \cdot \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{v} \cdot \boldsymbol{g}_{\rm Lorentz} + S_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \boldsymbol{f}_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{P}_{\rm cr} = - \qquad \frac{\boldsymbol{b}\boldsymbol{b}}{3\kappa_{\pm}} \cdot [\boldsymbol{f}_{\rm cr} - \boldsymbol{w}_{\pm}(\varepsilon_{\rm cr} + P_{\rm cr})] - \boldsymbol{g}_{\rm Lorentz} + \boldsymbol{S}_{f}$$

Alfvén wave velocity in lab frame: $w_{\pm} = v \pm v_a$, CR pressure tensor $P_{cr} = P_{cr} \mathbf{1}$, CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

• lab-frame equ's for radiation energy and momentum density, ε and f/c^2 (Mihalas & Mihalas 1984, Lowrie+ 1999):

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot \mathbf{f} = -\sigma_{s} \mathbf{v} \cdot [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_{a}$$
$$\frac{1}{c^{2}} \frac{\partial \mathbf{f}}{\partial t} + \nabla \cdot \mathbf{P} = -\sigma_{s} \quad [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_{a} \mathbf{v}$$

• **solution:** transform in comoving frame and project out gyrokinetics!

Alfvén-wave regulated CR transport

comoving equ's for CR energy and momentum density (along B), ε_{cr} and f_{cr}/c², and Alfvén-wave energy densities ε_{a,±} (Thomas & CP 2019)

$$\frac{\partial \varepsilon_{\rm cr}}{\partial t} + \boldsymbol{\nabla} \cdot [\boldsymbol{\nu}(\varepsilon_{\rm cr} + \boldsymbol{P}_{\rm cr}) + \boldsymbol{b}f_{\rm cr}] = \boldsymbol{\nu} \cdot \boldsymbol{\nabla}\boldsymbol{P}_{\rm cr} - \frac{\boldsymbol{v}_{\rm a}}{3\kappa_{+}} \left[f_{\rm cr} - \boldsymbol{v}_{\rm a}(\varepsilon_{\rm cr} + \boldsymbol{P}_{\rm cr})\right] + \frac{\boldsymbol{v}_{\rm a}}{3\kappa_{-}} \left[f_{\rm cr} + \boldsymbol{v}_{\rm a}(\varepsilon_{\rm cr} + \boldsymbol{P}_{\rm cr})\right],$$

$$\frac{\partial f_{\rm cr}/c^2}{\partial t} + \nabla \cdot \left(\boldsymbol{v} f_{\rm cr}/c^2 \right) + \boldsymbol{b} \cdot \nabla P_{\rm cr} = -(\boldsymbol{b} \cdot \nabla \boldsymbol{v}) \cdot (\boldsymbol{b} f_{\rm cr}/c^2) - \frac{1}{3\kappa_+} \left[f_{\rm cr} - v_{\rm a}(\varepsilon_{\rm cr} + P_{\rm cr}) \right] - \frac{1}{3\kappa_-} \left[f_{\rm cr} + v_{\rm a}(\varepsilon_{\rm cr} + P_{\rm cr}) \right],$$

$$\begin{split} \frac{\partial \varepsilon_{\mathrm{a},\pm}}{\partial t} + \boldsymbol{\nabla} \cdot \left[\boldsymbol{\nu} (\varepsilon_{\mathrm{a},\pm} + \boldsymbol{P}_{\mathrm{a},\pm}) \pm \boldsymbol{\nu}_{\mathrm{a}} \boldsymbol{b} \varepsilon_{\mathrm{a},\pm} \right] &= \boldsymbol{\nu} \cdot \boldsymbol{\nabla} \boldsymbol{P}_{\mathrm{a},\pm} \\ &\pm \frac{\boldsymbol{\nu}_{\mathrm{a}}}{3\kappa_{\pm}} \left[f_{\mathrm{cr}} \mp \boldsymbol{\nu}_{\mathrm{a}} (\varepsilon_{\mathrm{cr}} + \boldsymbol{P}_{\mathrm{cr}}) \right] - \boldsymbol{S}_{\mathrm{a},\pm}. \end{split}$$



Wave-particle interactions CR driven instabilities CR hydrodynamics

CR interactions with Alfvén waves



Wave-particle interactions CR driven instabilities CR hydrodynamics

CR interactions with Alfvén waves



Wave-particle interactions CR driven instabilities CR hydrodynamics

Non-equilibrium CR streaming and diffusion

Coupling the evolution of CR and Alfvén wave energy densities



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

١P

Wave-particle interactions CR driven instabilities CR hydrodynamics

Non-equilibrium CR streaming and diffusion

Varying damping rate of Alfvén waves modulates the diffusivity of solution



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

1. Galaxy simulations with cosmic ray feedback



Thomas, CP, Pakmor (2023) *Cosmic ray-driven galactic winds: transport modes of cosmic rays and Alfvén-wave dark regions*

MHD + Alfvén wave regulated CR hydrodynamics: $10^{11} M_{\odot}$ halo



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

CR interactions with Alfvén waves



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Wind launching



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

AIP

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Magnetic field topology



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Wind launching





ъ

Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds

Wind launching



Thomas, CP, Pakmor (2023)

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds

Wind launching



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

AIP

Cosmic ray driven winds

Wind properties



Temperature [K] 10^4 10^{5} 10^{6}



Thomas, CP, Pakmor (2023)







Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

AIP

What is the origin of the Alfvén wave dark regions?



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

What is the origin of the Alfvén wave dark regions?





CRs faster than AWs AWs gain energy

Cosmic Rays and Magnetic Fields in Galaxy Formation



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

What is the origin of the Alfvén wave dark regions?





CRs faster than AWs AWs gain energy



CRs slower than AWs AWs lose energy



Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Parallel CR diffusion coefficient

 $\begin{array}{c} {\rm CR\ energy\ per\ pixel\ [erg]}\\ 10^{47}\ 10^{48}\ 10^{49}\ 10^{50}\ 10^{51}\ 10^{52}\ 10^{53}\ 10^{54} \end{array}$



Thomas, CP, Pakmor (2023)

ъ

The CR diffusion coefficient is not constant but strongly depends on environment!



Introduction Cosmic ray Cosmic ray transport Galactic m Cosmic rays in galaxy formation Cosmologi

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Origin and growth of magnetic fields

The general picture:

• **Origin.** Magnetic fields are generated by 1. electric currents sourced by a phase transition in the early universe or 2. by the Biermann battery



< 合

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Origin and growth of magnetic fields

The general picture:

- **Origin.** Magnetic fields are generated by 1. electric currents sourced by a phase transition in the early universe or 2. by the Biermann battery
- Growth. A small-scale (fluctuating) dynamo is an MHD process, in which the kinetic (turbulent) energy is converted into magnetic energy: the mechanism relies on magnetic fields to become stronger when the field lines are stretched



AIP

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Origin and growth of magnetic fields

The general picture:

- **Origin.** Magnetic fields are generated by 1. electric currents sourced by a phase transition in the early universe or 2. by the Biermann battery
- Growth. A small-scale (fluctuating) dynamo is an MHD process, in which the kinetic (turbulent) energy is converted into magnetic energy: the mechanism relies on magnetic fields to become stronger when the field lines are stretched
- Saturation. Field growth stops at a sizeable fraction of the turbulent energy when magnetic forces become strong enough to resist the stretching and folding motions





Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

2. Galactic magnetic dynamo



CP, Werhahn, Pakmor, Girichidis, Simpson (2022)

Simulating radio synchrotron emission in star-forming galaxies: small-scale magnetic dynamo and the origin of the far-infrared-radio correlation

MHD + cosmic ray advection + diffusion: $\{10^{10}, 10^{11}, 3 \times 10^{11}, 10^{12}\} M_{\odot}$



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Time evolution of SFR and energy densities



 cosmic ray (CR) pressure feedback suppresses SFR more in smaller galaxies

- energy budget in disks is dominated by CR pressure
- magnetic growth faster in Milky Way galaxies than in dwarfs



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formatio

Identifying different growth phases



• 1st phase: adiabatic growth with $B \propto \rho^{2/3}$ (isotropic collapse)



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formatio

Identifying different growth phases



- 1st phase: adiabatic growth with $B \propto \rho^{2/3}$ (isotropic collapse)
- 2^{nd} phase: additional growth at high density ρ with small dynamical times $t_{dyn} \sim (G\rho)^{-1/2}$



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formatio

Identifying different growth phases



- 1st phase: adiabatic growth with $B \propto \rho^{2/3}$ (isotropic collapse)
- 2^{nd} phase: additional growth at high density ρ with small dynamical times $t_{dyn} \sim (G\rho)^{-1/2}$
- 3rd phase: growth migrates to lower ρ on larger scales $\propto \rho^{-1/3}$



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Studying growth rate with numerical resolution



CP+ (2022)

• faster magnetic growth in higher resolution simulations and larger halos, numerical convergence for $N \gtrsim 10^6$



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Studying growth rate with numerical resolution



- faster magnetic growth in higher resolution simulations and larger halos, numerical convergence for N ≥ 10⁶
- 1st phase: adiabatic growth (independent of resolution)



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formatior

Studying growth rate with numerical resolution



- faster magnetic growth in higher resolution simulations and larger halos, numerical convergence for N ≥ 10⁶
- 1st phase: adiabatic growth (independent of resolution)
- 2nd phase: small-scale dynamo with resolution-dep. growth rate

$$\Gamma = \frac{\mathscr{Y}}{\mathscr{L}} \operatorname{Re}_{\operatorname{num}}^{1/2}, \quad \operatorname{Re}_{\operatorname{num}} = \frac{\mathscr{L}\mathscr{Y}}{\mathcal{V}_{\operatorname{num}}} = \frac{3\mathscr{L}\mathscr{Y}}{d_{\operatorname{cell}} v_{\operatorname{th}}}$$

Cosmic ray driven winds Galactic magnetic dynamo

Exponential field growth in kinematic regime



 corrugated accretion shock dissipates kinetic energy from gravitational infall, injects vorticity that decays into turbulence, and drives a small-scale dynamo



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formatio

Dynamo saturation on small scales while λ_B increases



 supersonic velocity shear between the rotationally supported cool disk and hotter CGM: excitation of Kelvin-Helmholtz body modes that interact and drive a small-scale dynamo



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Kinetic and magnetic power spectra

Fluctuating small-scale dynamo in different analysis regions



- $E_B(k)$ superposition of form factor and turbulent spectrum
- pure turbulent spectrum outside steep central B profile



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

3. Cosmological galaxy formation

© Volker Springel/Max Planck for Astrophysics/Science Photo Library

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Cosmic rays in cosmological galaxy simulations

The galaxy formation model

- primordial and metal line cooling
- sub-resolution model for star formation (Springel+ 03)
- mass and metal return from stars to ISM
- cold dense gas stabilised by pressurised ISM
- thermal and kinetic energy from supernovae modelled by isotropic wind – launched outside of SF region
- black hole seeding and accretion model (Springel+ 05)
- thermal feedback from AGN in radio and quasar mode
- uniform magnetic field of 10⁻¹⁰ G seeded at z = 128

Simulation suite (Buck, CP+ 2020)

- 2 galaxies, baryons with 5 \times 10⁴ M_{\odot} \sim 5 \times 10⁶ resolution elements in halo, 2 \times 10⁶ star particles
- 4 models with different CR physics for each galaxy:
 - no CRs
 - CR advection
 - + CR anisotropic diffusion
 - + CR Alfvén wave cooling



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Cosmic rays in cosmological galaxy simulations

Auriga MHD models: CR transport changes disk sizes



Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Cosmic rays in cosmological galaxy simulations

Auriga MHD models: CR transport modifies the circum-galactic medium



Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

AIP

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Conclusions for cosmic ray physics in galaxies

CR hydrodynamics:

- moment expansion similar to radiation hydrodynamics
- novel theory of CR transport mediated by Alfvén waves and coupled to magneto-hydrodynamics
- Diffusion coefficient emerges from CR-wave interactions



Conclusions for cosmic ray physics in galaxies

CR hydrodynamics:

- moment expansion similar to radiation hydrodynamics
- novel theory of CR transport mediated by Alfvén waves and coupled to magneto-hydrodynamics
- Diffusion coefficient emerges from CR-wave interactions

CRs and magnetic fields in galaxy formation:

- small-scale dynamo grows magnetic field to equipartition with turbulent energy density
- CR feedback drives galactic winds & slows down star formation
- CRs modify galaxy disk sizes and the circumgalactic medium



A B + A B +
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

3 1 4 3

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

PICOGAL: From Flasma Kinetics to COsmological GALaxy Formation





Christoph Pfrommer

Cosmic Rays and Magnetic Fields in Galaxy Formation

Cosmic ray driven winds Galactic magnetic dynamo Cosmological galaxy formation

Literature for the talk

CR hydrodynamics:

- Pfrommer, Pakmor, Schaal, Simpson, Springel, *Simulating cosmic ray physics on a moving mesh*, 2017a, MNRAS, 465, 4500.
- Thomas & Pfrommer, Cosmic-ray hydrodynamics: Alfvén-wave regulated transport of cosmic rays, 2019, MNRAS, 485, 2977.
- Thomas, Pfrommer, Pakmor, A finite volume method for two-moment cosmic-ray hydrodynamics on a moving mesh, 2021, MNRAS, 503, 2242.
- Thomas & Pfrommer, Comparing different closure relations for cosmic ray hydrodynamics, 2022, MNRAS, 509, 4803.

Magnetic dynamos and CR feedback in galaxy formation:

- Pfrommer, Werhahn, Pakmor, Girichidis, Simpson, Simulating radio synchrotron emission in star-forming galaxies: small-scale magnetic dynamo and the origin of the far infrared-radio correlation, 2022, MNRAS, 515, 4229.
- Thomas, Pfrommer, Pakmor, Cosmic ray-driven galactic winds: transport modes of cosmic rays and Alfvén-wave dark regions, 2023, MNRAS, 521, 3023.



イロト イ理ト イヨト イヨト