



Cosmic rays in galaxy formation: instabilities, transport and feedback

Christoph Pfrommer¹

in collaboration with

PhD students: K. Ehler¹, R. Lemmerz¹, **T. Thomas¹**, **M. Werhahn¹**,
J. Whittingham¹, G. Winner¹

Postdocs: T. Berlok¹, T. Buck¹, P. Girichidis¹, **M. Shalaby¹**, M. Sparre^{2,1}
M. Pais³, E. Puchwein¹, R. Pakmor⁴, V. Springel⁴, T. Enßlin⁴, C. Simpson⁵

¹AIP Potsdam, ²U of Potsdam, ³Hebrew U, ⁴MPA Garching, ⁵U of Chicago

MIAPP, Munich, July 2021

Outline

1 Cosmic ray driven instabilities

- Introduction
- Intermediate instability
- Overview and applications

2 Cosmic ray transport

- CR propagation
- CR hydrodynamics
- Radio synchrotron harps

3 Galaxy formation

- Cosmic ray driven winds
- Non-thermal emission
- Conclusions



Does plasma physics matter in galaxy formation?



Does plasma physics matter in galaxy formation?

Can (sub-)galactic observations teach us plasma physics?

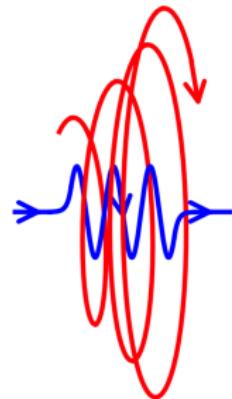


CR streaming instability

- **Cosmic ray (CR) streaming instability:**

Kulsrud & Pearce 1969

- if $v_{\text{cr}} > v_a$, CR flux excites and amplifies an Alfvén wave field in resonance with the gyroradii of CRs
- scattering off of this wave field limits the (GeV) CRs' bulk speed $\sim v_a$
- wave damping: transfer of CR energy and momentum to the thermal gas

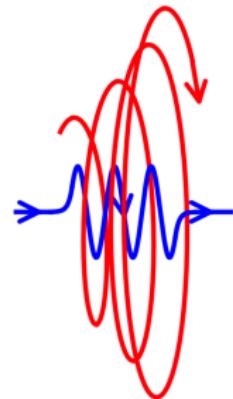


CR streaming instability

- **Cosmic ray (CR) streaming instability:**

Kulsrud & Pearce 1969

- if $v_{\text{cr}} > v_a$, CR flux excites and amplifies an Alfvén wave field in resonance with the gyroradii of CRs
- scattering off of this wave field limits the (GeV) CRs' bulk speed $\sim v_a$
- wave damping: **transfer of CR energy and momentum to the thermal gas**



→ CRs exert pressure on thermal gas via scattering on Alfvén waves

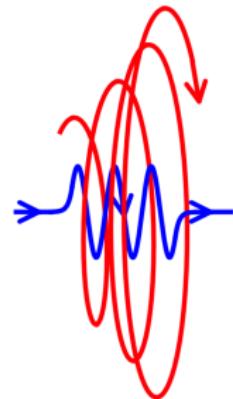


CR streaming instability

- **Cosmic ray (CR) streaming instability:**

Kulsrud & Pearce 1969

- if $v_{\text{cr}} > v_a$, CR flux excites and amplifies an Alfvén wave field in resonance with the gyroradii of CRs
- scattering off of this wave field limits the (GeV) CRs' bulk speed $\sim v_a$
- wave damping: transfer of CR energy and momentum to the thermal gas



→ CRs exert pressure on thermal gas via scattering on Alfvén waves

weak wave damping: strong coupling → CR stream with waves

strong wave damping: less waves to scatter → CR diffusion prevails



Coupling of CRs to the background plasma

- **extrinsic confinement:** scattering off of turbulence cascaded from large scales via supernovae, stellar winds, etc.
⇒ important for confinement of TeV CRs



Coupling of CRs to the background plasma

- **extrinsic confinement:** scattering off of turbulence cascaded from large scales via supernovae, stellar winds, etc.
⇒ important for confinement of TeV CRs
- **intrinsic confinement:** CRs drive unstable plasma wave modes (e.g., Alfvén waves), and then scatter off of them
⇒ most important mechanism for GeV CR confinement



Coupling of CRs to the background plasma

- **extrinsic confinement:** scattering off of turbulence cascaded from large scales via supernovae, stellar winds, etc.
 \Rightarrow important for confinement of TeV CRs
- **intrinsic confinement:** CRs drive unstable plasma wave modes (e.g., Alfvén waves), and then scatter off of them
 \Rightarrow most important mechanism for GeV CR confinement
- **dispersion relation** ($\Omega_{e,0} = -m_i/m_e \times \Omega_{i,0}$, $\alpha = n_{\text{cr}}/n_i$): gyrotropic CR ion + electron beam propagates in background plasma

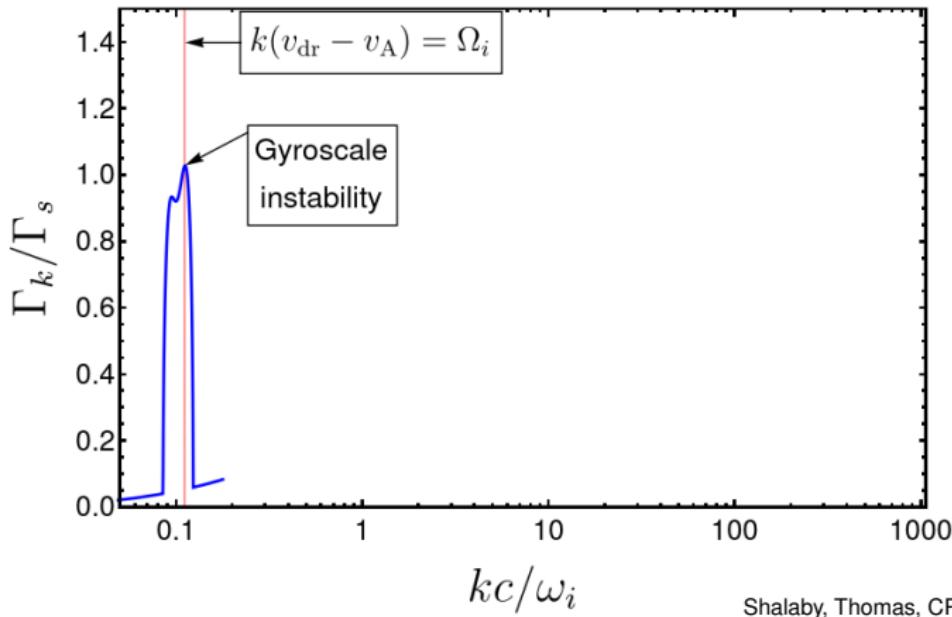
$$\frac{k^2 c^2}{\omega^2} - 1 = \frac{\omega_i^2}{\omega(-\omega \pm \Omega_{i,0})} + \frac{\omega_e^2}{\omega(-\omega \pm \Omega_{e,0})} \quad \leftarrow \text{background}$$

$$\text{CRe} \Rightarrow + \frac{\alpha \omega_e^2}{\gamma_e \omega^2} \left\{ \frac{\omega - kv_{\text{dr}}}{kv_{\text{dr}} - \omega \mp \Omega_{e,0}/\gamma_e} \right\}$$

$$\text{CRI} \Rightarrow + \frac{\alpha \omega_i^2}{\gamma_i \omega^2} \left\{ \frac{\omega - kv_{\text{dr}}}{kv_{\text{dr}} - \omega \pm \Omega_i} - \frac{v_{\perp}^2 (k^2 c^2 - \omega^2) / c^2}{2 (kv_{\text{dr}} - \omega \pm \Omega_i)^2} \right\}$$



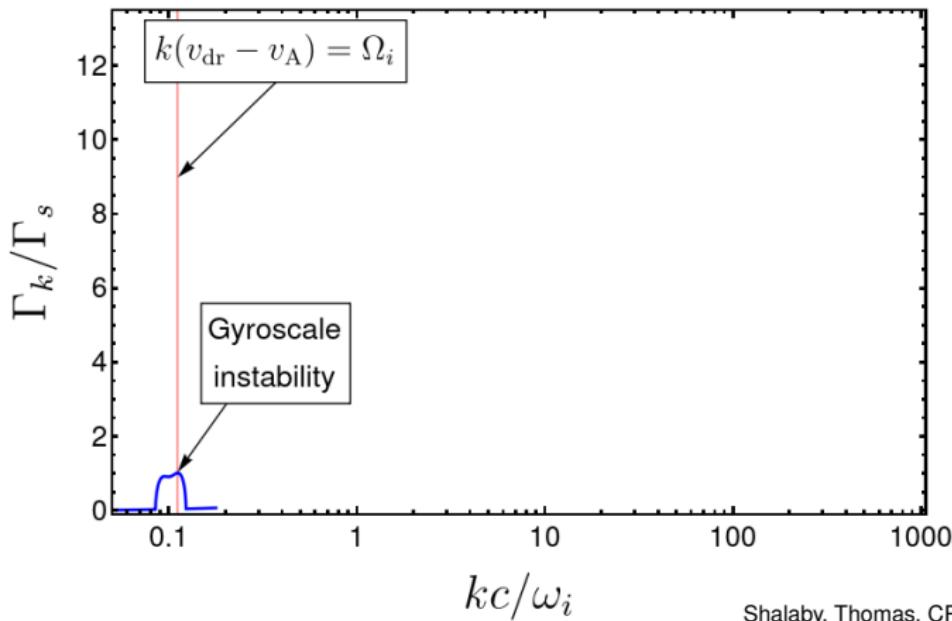
CR driven instabilities – growth rates



Shalaby, Thomas, CP (2021)

- gyro-resonant instability of gyrotropic CR population

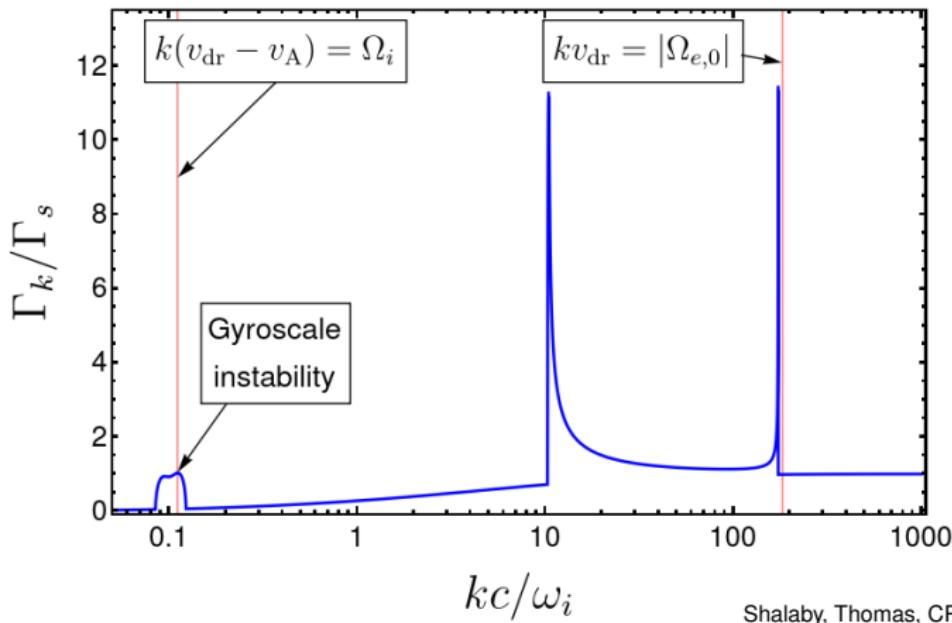
CR driven instabilities – growth rates



Shalaby, Thomas, CP (2021)

- gyro-resonant instability of gyrotropic CR population

CR driven instabilities – growth rates

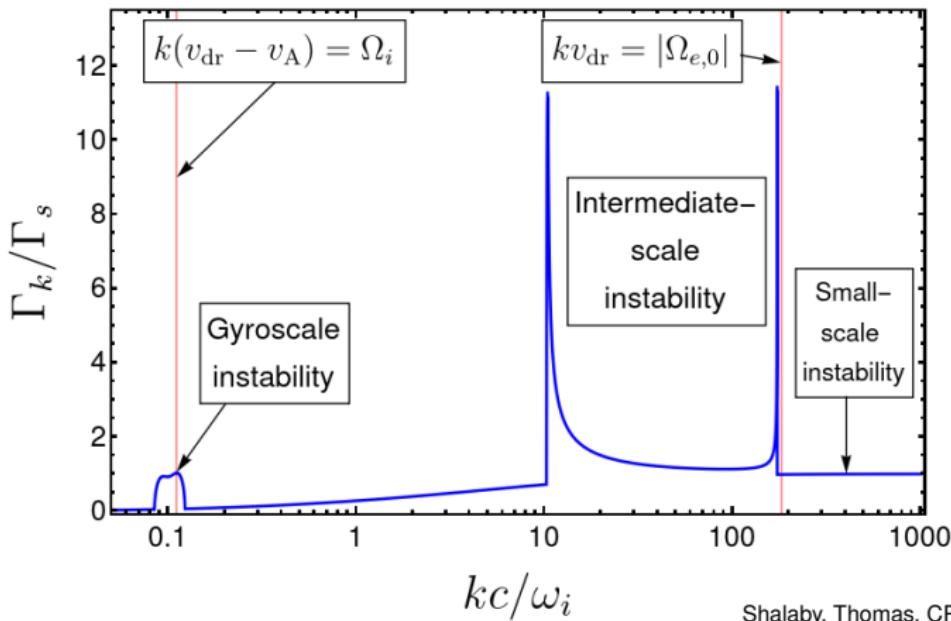


Shalaby, Thomas, CP (2021)



- new intermediate-scale instability of gyrotrropic CR population

CR driven instabilities – growth rates

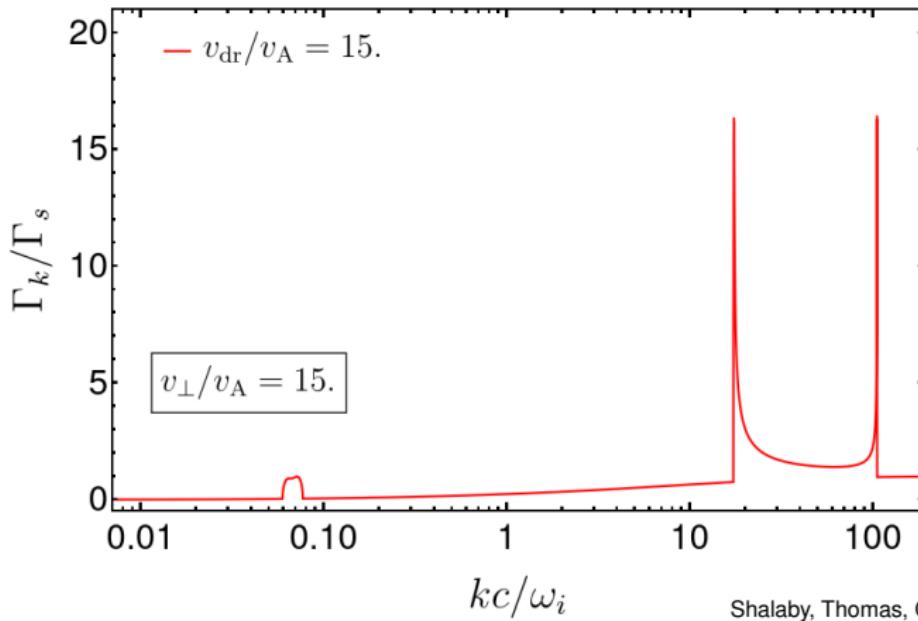


Shalaby, Thomas, CP (2021)



- new intermediate-scale instability of gyrotrropic CR population

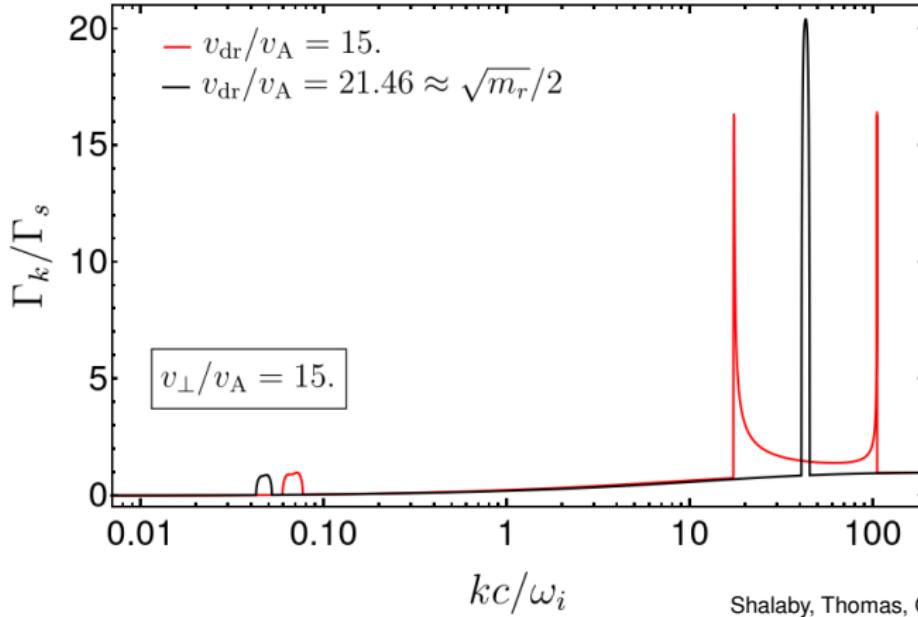
CR driven intermediate-scale instability



Shalaby, Thomas, CP (2021)

- low CR drift speed: two instability peaks

CR driven intermediate-scale instability

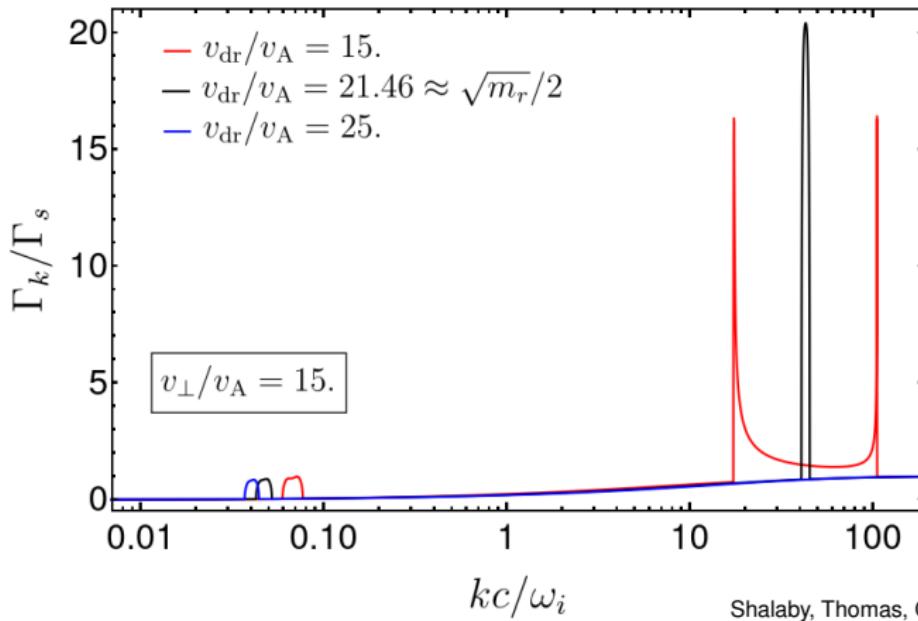


Shalaby, Thomas, CP (2021)



- for CR drift speed $v_{\text{dr}} \approx \sqrt{\frac{m_i}{m_e}} \frac{v_A}{2}$: two instability peaks merge

CR driven intermediate-scale instability

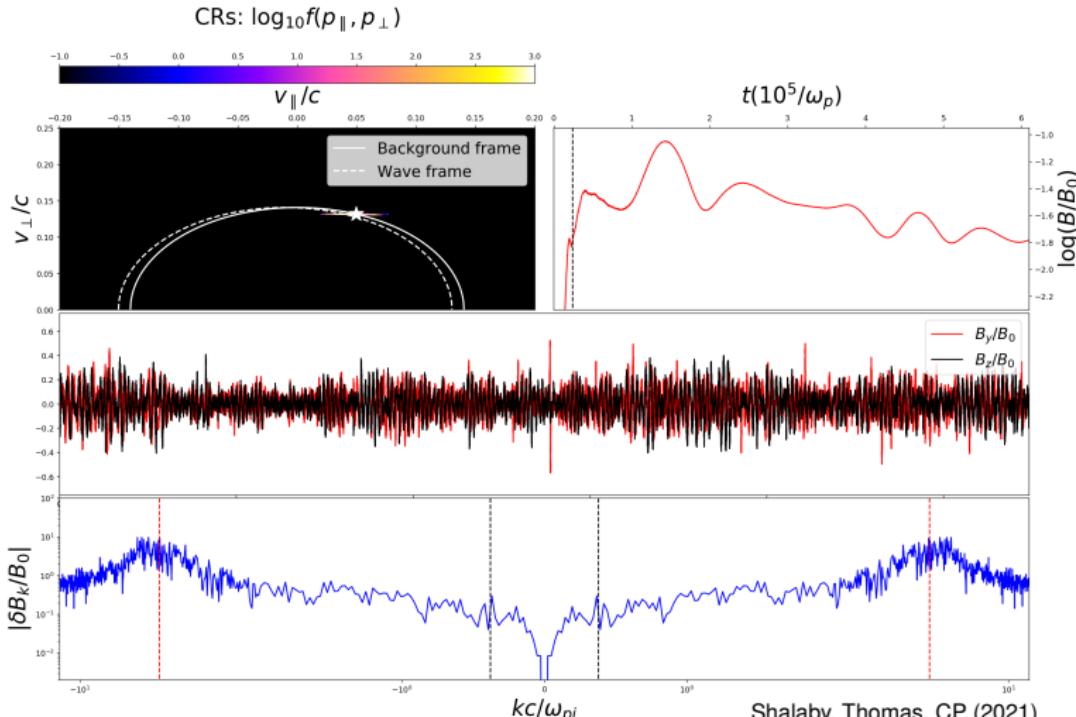


Shalaby, Thomas, CP (2021)

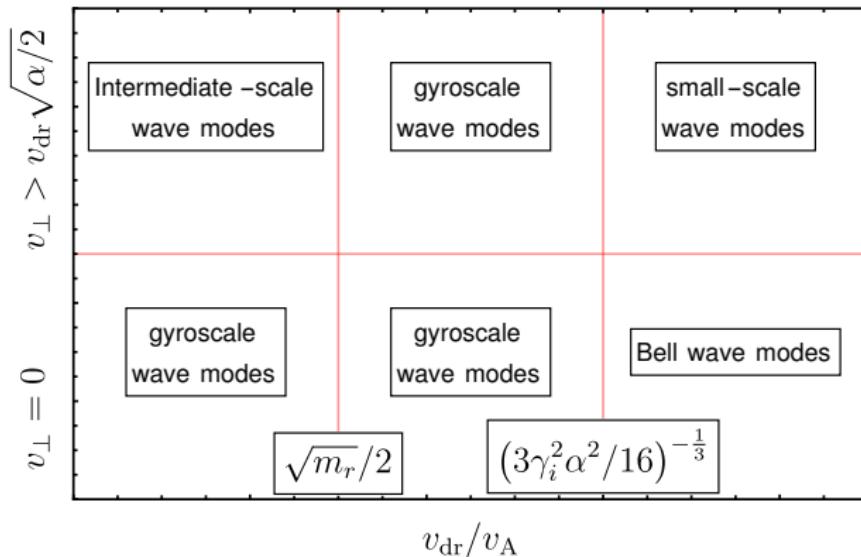
- for $v_{dr} > \sqrt{\frac{m_i}{m_e}} \frac{v_A}{2}$: intermediate-scale instability quenched

Cosmic ray driven instabilities

Growth of the intermediate-scale and the gyro-resonant instability



Regimes of CR driven instabilities



Shalaby, Thomas, CP (2021)

- where $\alpha = \frac{n_{\text{cr}}}{n_i}$ is the CR number fraction, $m_r = \frac{m_i}{m_e}$ is the mass ratio, and γ_i is the Lorentz factor of CR ions



The intermediate-scale instability

Properties of the intermediate-scale instability:

- growth rate $\Gamma_{\text{inter}} \gg \Gamma_{\text{gyro}}$ and excites broad spectral support
- unstable modes are background ion-cyclotron waves in the comoving CR frame
- condition for growth:
$$\frac{v_{\text{dr}}}{v_A} < \frac{1}{2} \sqrt{\frac{m_i}{m_e}}$$



AIP

The intermediate-scale instability

Properties of the intermediate-scale instability:

- growth rate $\Gamma_{\text{inter}} \gg \Gamma_{\text{gyro}}$ and excites broad spectral support
- unstable modes are background ion-cyclotron waves in the comoving CR frame
- condition for growth: $\frac{V_{\text{dr}}}{V_A} < \frac{1}{2} \sqrt{\frac{m_i}{m_e}}$

Possible implications of this new instability:

- couples CRs more tightly to background plasma and strengthens CR feedback in galaxies and galaxy clusters
- enables electron heating at shocks and injection into diffusive shock acceleration
- slows down CR escape from the sites of particle acceleration
→ brighter gamma-ray halos

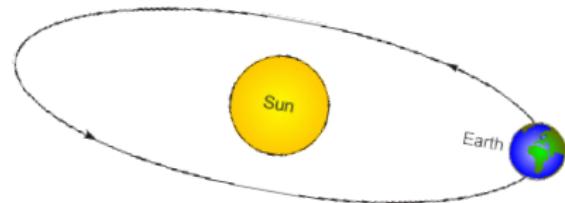


Cosmic ray transport: an extreme multi-scale problem



Milky Way-like galaxy:

$$r_{\text{gal}} \sim 10^4 \text{ pc}$$



gyro-orbit of GeV cosmic ray:

$$r_{\text{cr}} = \frac{p_{\perp}}{e B_{\mu G}} \sim 10^{-6} \text{ pc} \sim \frac{1}{4} \text{ AU}$$

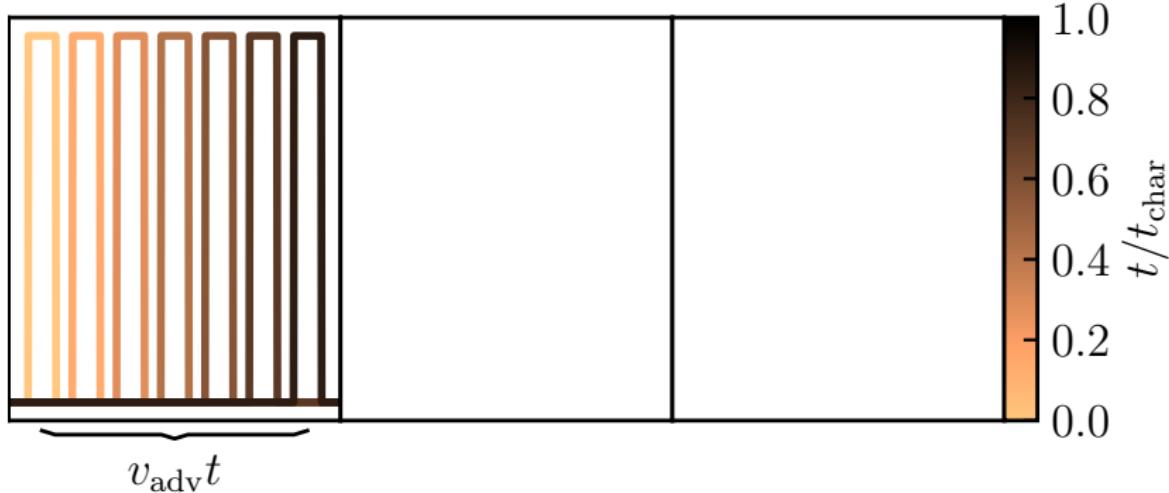
⇒ need to develop a **fluid theory for a collisionless, non-Maxwellian component!**

Zweibel (2017), Jiang & Oh (2018), Thomas & CP (2019)



Modes of CR propagation

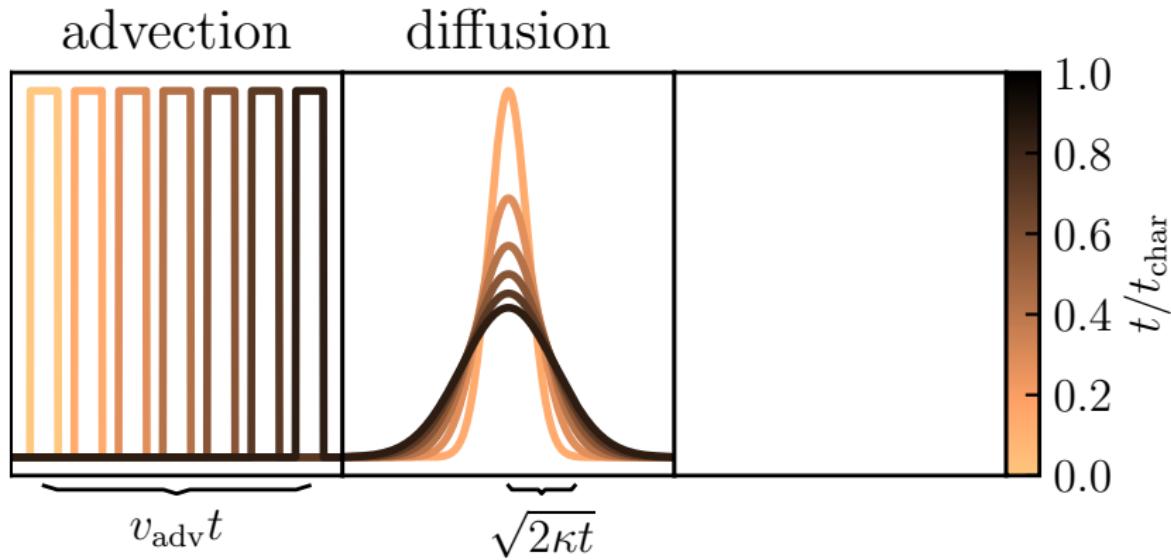
advection



Thomas, CP, Enßlin (2020)



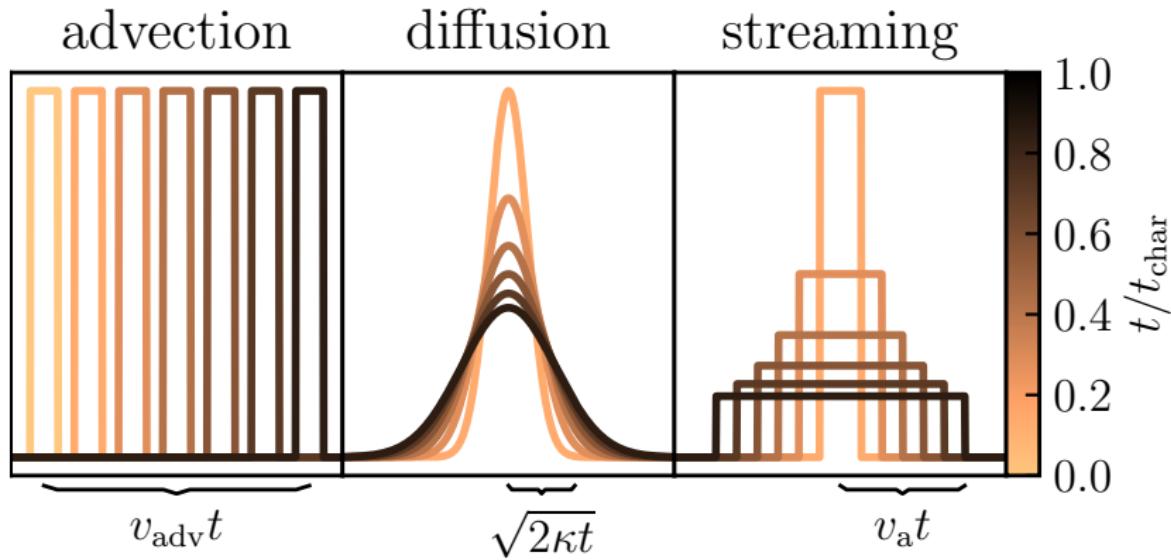
Modes of CR propagation



Thomas, CP, Enßlin (2020)



Modes of CR propagation



Thomas, CP, Enßlin (2020)



1-moment CR hydrodynamics (steady-state flux)

- total CR velocity $\mathbf{v}_{\text{cr}} = \mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}}$ (where $\mathbf{v} \equiv \mathbf{v}_{\text{gas}}$)



1-moment CR hydrodynamics (steady-state flux)

- total CR velocity $\mathbf{v}_{\text{cr}} = \mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}}$ (where $\mathbf{v} \equiv \mathbf{v}_{\text{gas}}$)
- CRs stream down their own pressure gradient relative to the gas

$$\mathbf{v}_{\text{st}} = \mathbf{v}_A \frac{\bar{\nu}_+ - \bar{\nu}_-}{\bar{\nu}_+ + \bar{\nu}_-},$$



1-moment CR hydrodynamics (steady-state flux)

- total CR velocity $\mathbf{v}_{\text{cr}} = \mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}}$ (where $\mathbf{v} \equiv \mathbf{v}_{\text{gas}}$)
- CRs stream** down their own pressure gradient relative to the gas,
CRs diffuse in the wave frame due to pitch angle scattering by MHD waves (both transports are along the local direction of \mathbf{B}):

$$\mathbf{v}_{\text{st}} = \mathbf{v}_A \frac{\bar{\nu}_+ - \bar{\nu}_-}{\bar{\nu}_+ + \bar{\nu}_-}, \quad \mathbf{v}_{\text{di}} = -\kappa_{\text{di}} \mathbf{b} \frac{\mathbf{b} \cdot \nabla \varepsilon_{\text{cr}}}{\varepsilon_{\text{cr}}}, \quad \kappa_{\text{di}} = \frac{c^2}{3(\bar{\nu}_+ + \bar{\nu}_-)}$$



1-moment CR hydrodynamics (steady-state flux)

- total CR velocity $\mathbf{v}_{\text{cr}} = \mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}}$ (where $\mathbf{v} \equiv \mathbf{v}_{\text{gas}}$)
- CRs stream** down their own pressure gradient relative to the gas,
CRs diffuse in the wave frame due to pitch angle scattering by MHD waves (both transports are along the local direction of \mathbf{B}):

$$\mathbf{v}_{\text{st}} = \mathbf{v}_A \frac{\bar{\nu}_+ - \bar{\nu}_-}{\bar{\nu}_+ + \bar{\nu}_-}, \quad \mathbf{v}_{\text{di}} = -\kappa_{\text{di}} \mathbf{B} \frac{\mathbf{b} \cdot \nabla \varepsilon_{\text{cr}}}{\varepsilon_{\text{cr}}}, \quad \kappa_{\text{di}} = \frac{c^2}{3(\bar{\nu}_+ + \bar{\nu}_-)}$$

- energy equations with $\varepsilon = \varepsilon_{\text{th}} + \rho v^2/2$:

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot [(\varepsilon + P_{\text{th}}) \mathbf{v}] = 0$$



1-moment CR hydrodynamics (steady-state flux)

- total CR velocity $\mathbf{v}_{\text{cr}} = \mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}}$ (where $\mathbf{v} \equiv \mathbf{v}_{\text{gas}}$)
- CRs stream** down their own pressure gradient relative to the gas,
CRs diffuse in the wave frame due to pitch angle scattering by MHD waves (both transports are along the local direction of \mathbf{B}):

$$\mathbf{v}_{\text{st}} = \mathbf{v}_A \frac{\bar{\nu}_+ - \bar{\nu}_-}{\bar{\nu}_+ + \bar{\nu}_-}, \quad \mathbf{v}_{\text{di}} = -\kappa_{\text{di}} \mathbf{b} \frac{\mathbf{b} \cdot \nabla \varepsilon_{\text{cr}}}{\varepsilon_{\text{cr}}}, \quad \kappa_{\text{di}} = \frac{c^2}{3(\bar{\nu}_+ + \bar{\nu}_-)}$$

- energy equations with $\varepsilon = \varepsilon_{\text{th}} + \rho v^2/2$:

$$\begin{aligned}\frac{\partial \varepsilon}{\partial t} + \nabla \cdot [(\varepsilon + P_{\text{th}} + P_{\text{cr}}) \mathbf{v}] &= P_{\text{cr}} \nabla \cdot \mathbf{v} - \mathbf{v}_{\text{st}} \cdot \nabla P_{\text{cr}} \\ \frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot [P_{\text{cr}} \mathbf{v}_{\text{st}} + \varepsilon_{\text{cr}} (\mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}})] &= -P_{\text{cr}} \nabla \cdot \mathbf{v} + \mathbf{v}_{\text{st}} \cdot \nabla P_{\text{cr}}\end{aligned}$$



1-moment CR hydrodynamics (steady-state flux)

- total CR velocity $\mathbf{v}_{\text{cr}} = \mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}}$ (where $\mathbf{v} \equiv \mathbf{v}_{\text{gas}}$)
- CRs stream** down their own pressure gradient relative to the gas,
CRs diffuse in the wave frame due to pitch angle scattering by MHD waves (both transports are along the local direction of \mathbf{B}):

$$\mathbf{v}_{\text{st}} = \mathbf{v}_A \frac{\bar{\nu}_+ - \bar{\nu}_-}{\bar{\nu}_+ + \bar{\nu}_-}, \quad \mathbf{v}_{\text{di}} = -\kappa_{\text{di}} \mathbf{b} \frac{\mathbf{b} \cdot \nabla \varepsilon_{\text{cr}}}{\varepsilon_{\text{cr}}}, \quad \kappa_{\text{di}} = \frac{c^2}{3(\bar{\nu}_+ + \bar{\nu}_-)}$$

- energy equations with $\varepsilon = \varepsilon_{\text{th}} + \rho v^2/2$:

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot [(\varepsilon + P_{\text{th}} + P_{\text{cr}}) \mathbf{v}] = P_{\text{cr}} \nabla \cdot \mathbf{v} - \mathbf{v}_{\text{st}} \cdot \nabla P_{\text{cr}}$$

$$\frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot [P_{\text{cr}} \mathbf{v}_{\text{st}} + \varepsilon_{\text{cr}} (\mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}})] = -P_{\text{cr}} \nabla \cdot \mathbf{v} + \mathbf{v}_{\text{st}} \cdot \nabla P_{\text{cr}}$$

$$\Leftrightarrow \frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot [\varepsilon_{\text{cr}} (\mathbf{v} + \mathbf{v}_{\text{st}} + \mathbf{v}_{\text{di}})] = -P_{\text{cr}} \nabla \cdot (\mathbf{v} + \mathbf{v}_{\text{st}})$$



AIP

2-moment CR vs. radiation hydrodynamics

- capitalize on **analogies of CR and radiation hydrodynamics** (Jiang & Oh 2018)
derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)



2-moment CR vs. radiation hydrodynamics

- capitalize on **analogies of CR and radiation hydrodynamics** (Jiang & Oh 2018)
 derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for **CR energy and momentum density, ε_{cr} and $\mathbf{f}_{\text{cr}}/c^2$**

$$\frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{f}_{\text{cr}} = -\mathbf{w}_{\pm} \cdot \frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{v} \cdot \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{P}_{\text{cr}} = -\frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_f$$

Alfvén wave velocity in lab frame: $\mathbf{w}_{\pm} = \mathbf{v} \pm \mathbf{v}_a$,
 CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$



2-moment CR vs. radiation hydrodynamics

- capitalize on **analogies of CR and radiation hydrodynamics** (Jiang & Oh 2018)
 derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for **CR energy and momentum density, ε_{cr} and $\mathbf{f}_{\text{cr}}/c^2$**

$$\frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{f}_{\text{cr}} = -\mathbf{w}_{\pm} \cdot \frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{v} \cdot \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{P}_{\text{cr}} = -\frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_f$$

Alfvén wave velocity in lab frame: $\mathbf{w}_{\pm} = \mathbf{v} \pm \mathbf{v}_a$,

CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

- lab-frame equ's for **radiation energy and momentum density, ε and \mathbf{f}/c^2**

(Mihalas & Mihalas, 1984, Lowrie+ 1999):

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot \mathbf{f} = -\sigma_s \mathbf{v} \cdot [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_a$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}}{\partial t} + \nabla \cdot \mathbf{P} = -\sigma_s [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_a \mathbf{v}$$



2-moment CR vs. radiation hydrodynamics

- capitalize on **analogies of CR and radiation hydrodynamics** (Jiang & Oh 2018)
 derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for **CR energy and momentum density, ε_{cr} and $\mathbf{f}_{\text{cr}}/c^2$**

$$\frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{f}_{\text{cr}} = -\mathbf{w}_{\pm} \cdot \frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{v} \cdot \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{P}_{\text{cr}} = -\frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_f$$

Alfvén wave velocity in lab frame: $\mathbf{w}_{\pm} = \mathbf{v} \pm \mathbf{v}_a$,

CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

- lab-frame equ's for **radiation energy and momentum density, ε and \mathbf{f}/c^2**
 (Mihalas & Mihalas, 1984, Lowrie+ 1999):

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot \mathbf{f} = -\sigma_s \mathbf{v} \cdot [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_a$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}}{\partial t} + \nabla \cdot \mathbf{P} = -\sigma_s [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_a \mathbf{v}$$

- **problem:** CR lab-frame equation requires resolving rapid gyrokinetics!



2-moment CR vs. radiation hydrodynamics

- capitalize on **analogies of CR and radiation hydrodynamics** (Jiang & Oh 2018)
 derive two-moment equations from CR Vlasov equation (Thomas & CP 2019)
- lab-frame equ's for **CR energy and momentum density, ε_{cr} and $\mathbf{f}_{\text{cr}}/c^2$**

$$\frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{f}_{\text{cr}} = -\mathbf{w}_{\pm} \cdot \frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{v} \cdot \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_{\varepsilon}$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}_{\text{cr}}}{\partial t} + \nabla \cdot \mathbf{P}_{\text{cr}} = -\frac{\mathbf{b}\mathbf{b}}{3\kappa_{\pm}} \cdot [\mathbf{f}_{\text{cr}} - \mathbf{w}_{\pm}(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \mathbf{g}_{\text{Lorentz}} + \mathbf{S}_f$$

Alfvén wave velocity in lab frame: $\mathbf{w}_{\pm} = \mathbf{v} \pm \mathbf{v}_a$,

CR scattering frequency $\bar{\nu}_{\pm} = c^2/(3\kappa_{\pm})$

- lab-frame equ's for **radiation energy and momentum density, ε and \mathbf{f}/c^2**
 (Mihalas & Mihalas, 1984, Lowrie+ 1999):

$$\frac{\partial \varepsilon}{\partial t} + \nabla \cdot \mathbf{f} = -\sigma_s \mathbf{v} \cdot [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_a$$

$$\frac{1}{c^2} \frac{\partial \mathbf{f}}{\partial t} + \nabla \cdot \mathbf{P} = -\sigma_s [\mathbf{f} - \mathbf{v} \cdot (\varepsilon \mathbf{1} + \mathbf{P})] + S_a \mathbf{v}$$

- **solution:** transform in comoving frame and project out gyrokinetics!



Alfvén-wave regulated CR transport

- comoving equ's for CR energy and momentum density (along \mathbf{B}), ε_{cr} and f_{cr}/c^2 , and Alfvén-wave energy densities $\varepsilon_{a,\pm}$ (Thomas & CP 2019)

$$\frac{\partial \varepsilon_{\text{cr}}}{\partial t} + \nabla \cdot [\mathbf{v}(\varepsilon_{\text{cr}} + P_{\text{cr}}) + \mathbf{b}f_{\text{cr}}] = \mathbf{v} \cdot \nabla P_{\text{cr}} - \frac{V_a}{3\kappa_+} [f_{\text{cr}} - V_a(\varepsilon_{\text{cr}} + P_{\text{cr}})] + \frac{V_a}{3\kappa_-} [f_{\text{cr}} + V_a(\varepsilon_{\text{cr}} + P_{\text{cr}})],$$

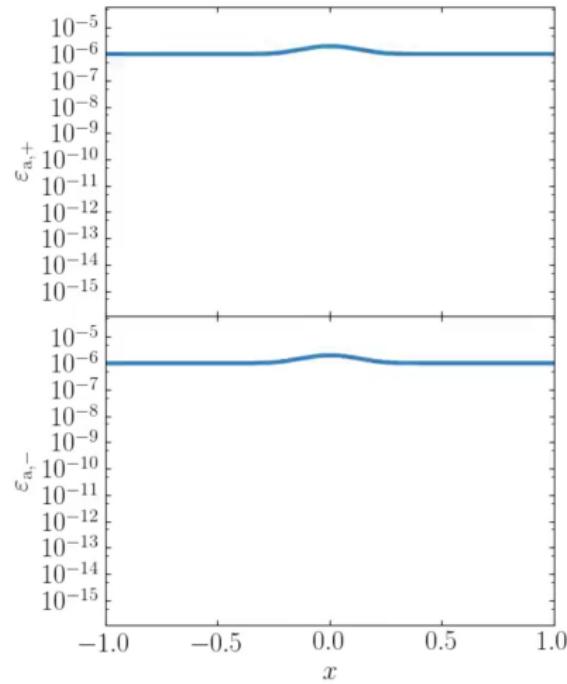
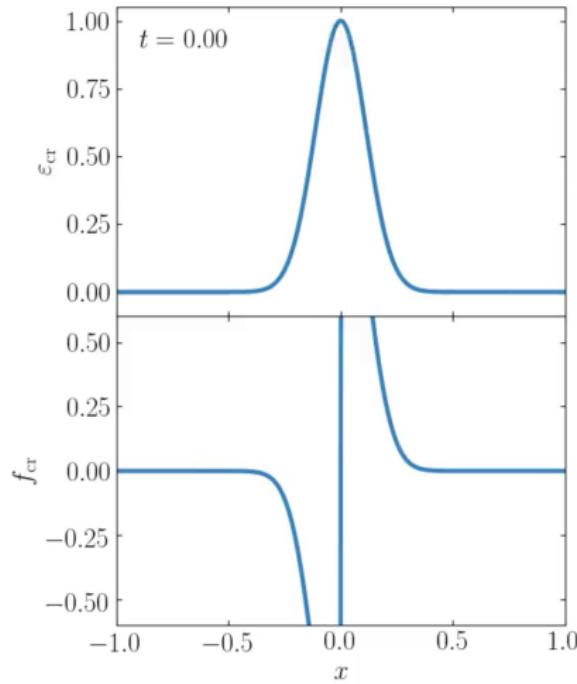
$$\frac{\partial f_{\text{cr}}/c^2}{\partial t} + \nabla \cdot \left(\mathbf{v}f_{\text{cr}}/c^2 \right) + \mathbf{b} \cdot \nabla P_{\text{cr}} = -(\mathbf{b} \cdot \nabla \mathbf{v}) \cdot (\mathbf{b}f_{\text{cr}}/c^2) - \frac{1}{3\kappa_+} [f_{\text{cr}} - V_a(\varepsilon_{\text{cr}} + P_{\text{cr}})] - \frac{1}{3\kappa_-} [f_{\text{cr}} + V_a(\varepsilon_{\text{cr}} + P_{\text{cr}})],$$

$$\frac{\partial \varepsilon_{a,\pm}}{\partial t} + \nabla \cdot [\mathbf{v}(\varepsilon_{a,\pm} + P_{a,\pm}) \pm V_a \mathbf{b} \varepsilon_{a,\pm}] = \mathbf{v} \cdot \nabla P_{a,\pm} \pm \frac{V_a}{3\kappa_\pm} [f_{\text{cr}} \mp V_a(\varepsilon_{\text{cr}} + P_{\text{cr}})] - S_{a,\pm}.$$



Non-equilibrium CR streaming and diffusion

Coupling the evolution of CR and Alfvén wave energy densities

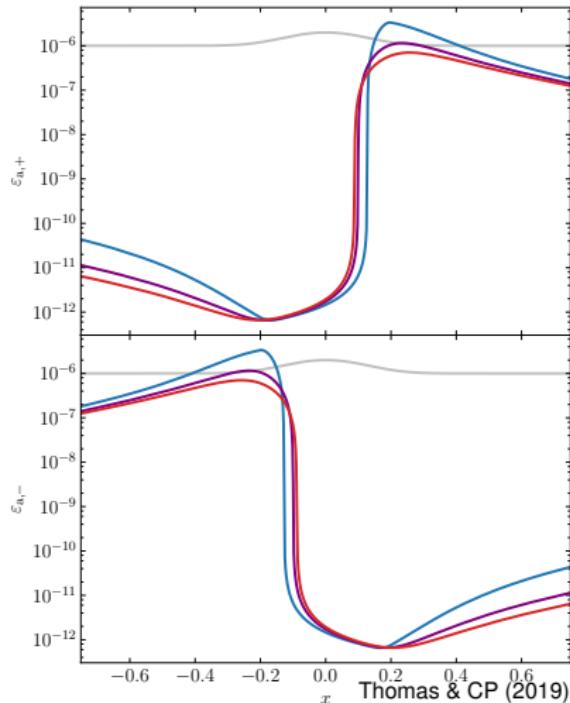
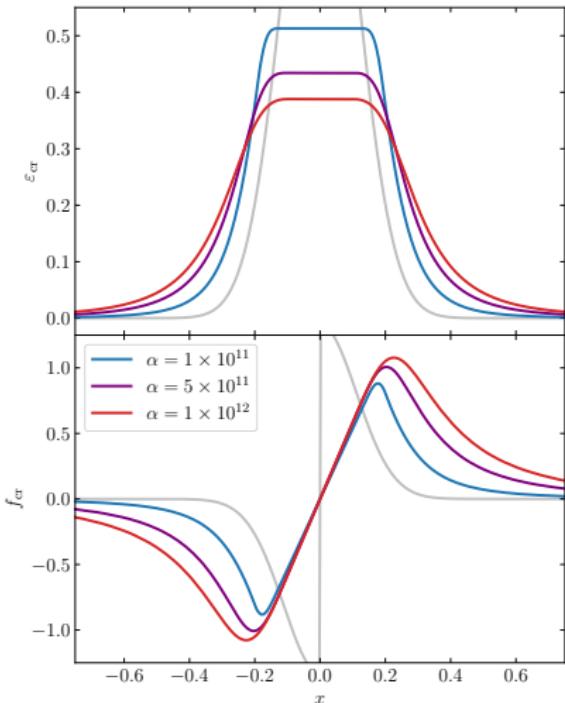


Thomas & CP (2019)



Non-equilibrium CR streaming and diffusion

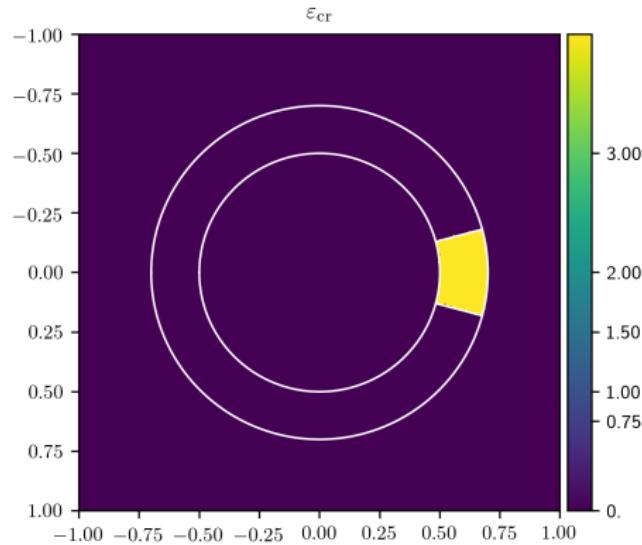
Varying damping rate of Alfvén waves modulates the diffusivity of solution



Anisotropic CR streaming and diffusion – AREPO

CR transport mediated by Alfvén waves and coupled to magneto-hydrodynamics

- CR streaming and diffusion along magnetic field lines in the self-confinement picture
- moment expansion similar to radiation hydrodynamics
- accounts for kinetic physics: non-linear Landau damping, gyro-resonant instability, ...
- Galilean invariant and causal transport
- energy and momentum conserving



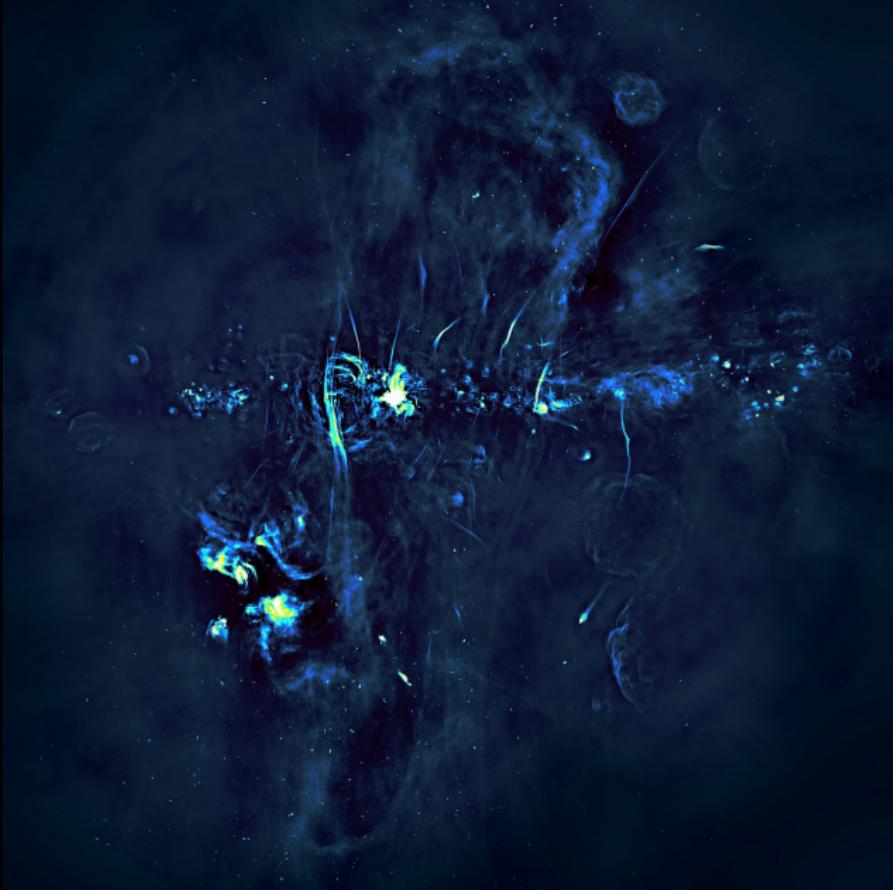
Thomas, CP, Pakmor (2021), Thomas & CP (2021)



AIP

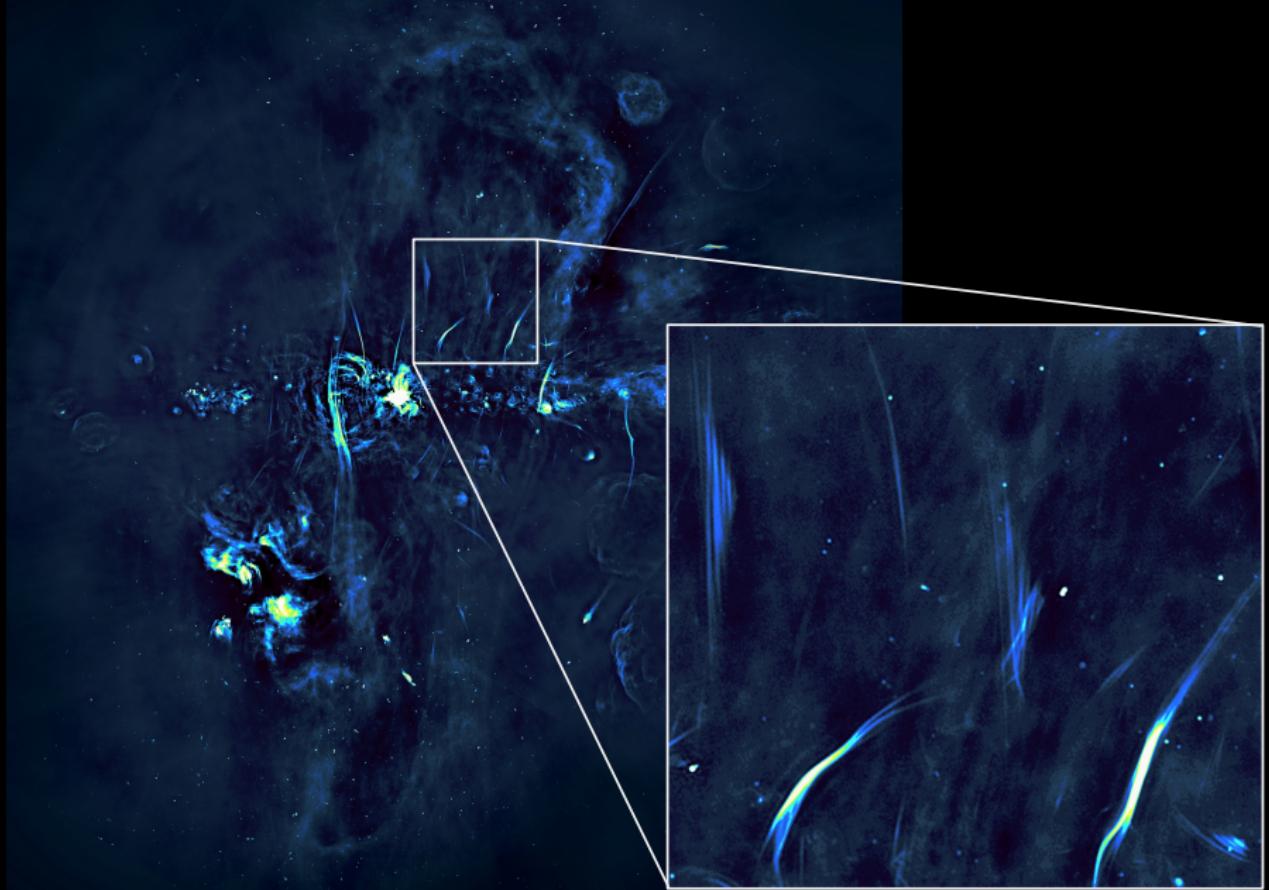
MeerKAT image of the Galactic Center

Haywood+ (Nature, 2019)



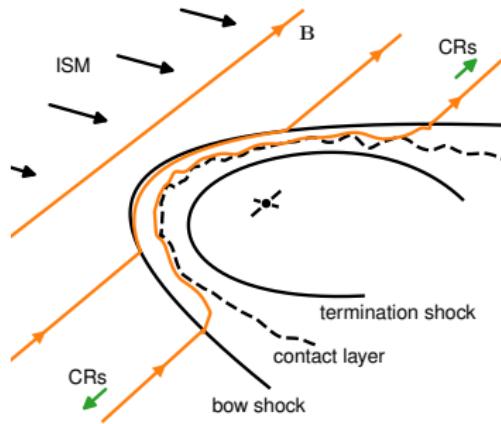
MeerKAT image of the Galactic Center

Haywood+ (Nature, 2019)



Radio synchrotron harps: the model

shock acceleration scenario

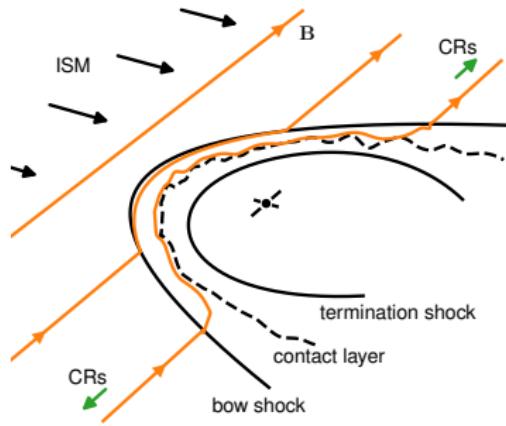


Thomas, CP, Enßlin (2020)

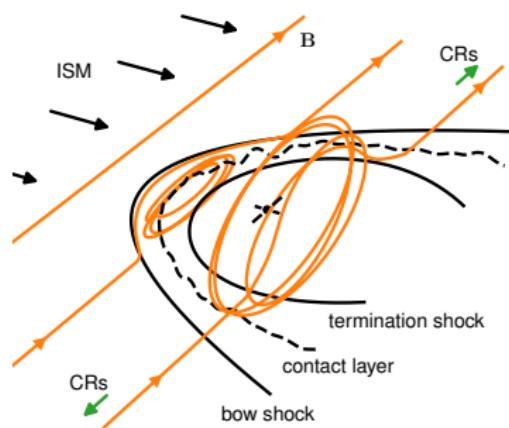


Radio synchrotron harps: the model

shock acceleration scenario



magnetic reconnection at pulsar wind

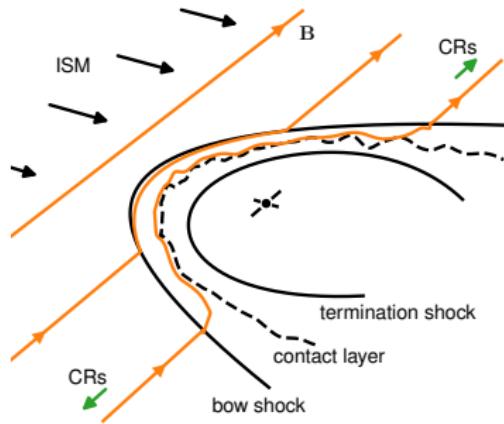


Thomas, CP, Enßlin (2020)



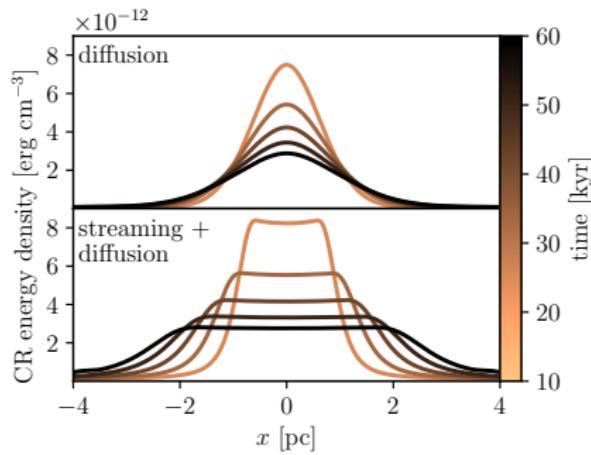
Radio synchrotron harps: the model

shock acceleration scenario

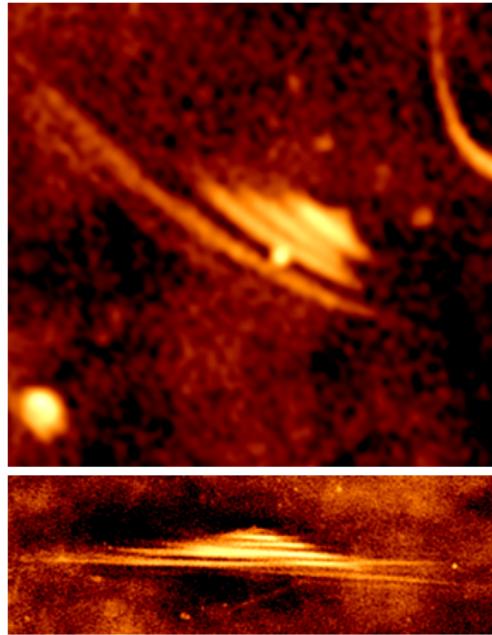


Thomas, CP, Enßlin (2020)

CR diffusion vs. streaming + diffusion



Radio synchrotron harps: testing CR propagation

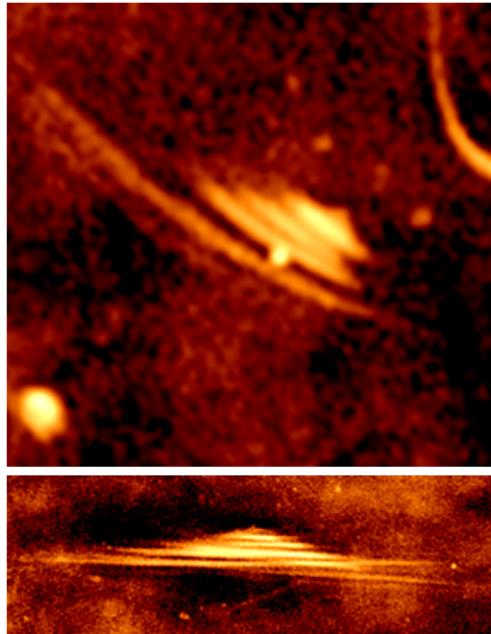


Haywood+ (Nature, 2019)



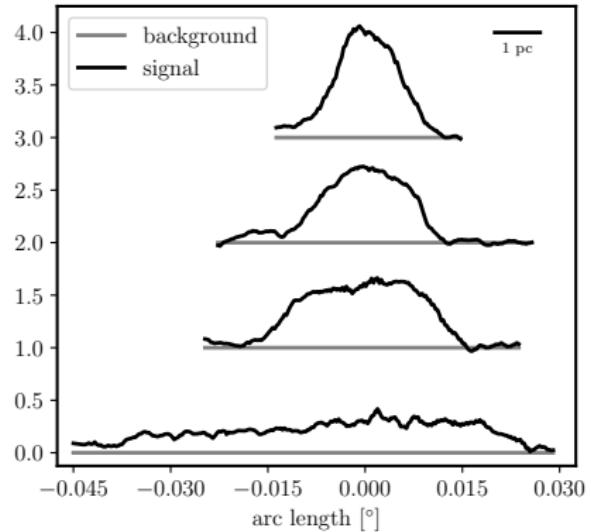
AIP

Radio synchrotron harps: testing CR propagation



Haywood+ (Nature, 2019)

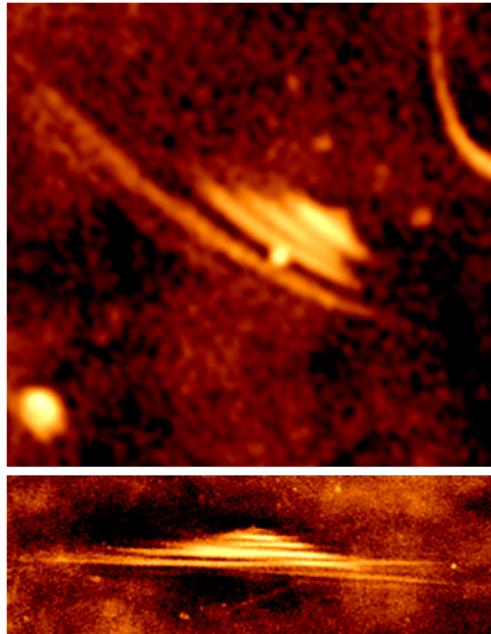
lateral radio profiles



Thomas, CP, Enßlin (2020)

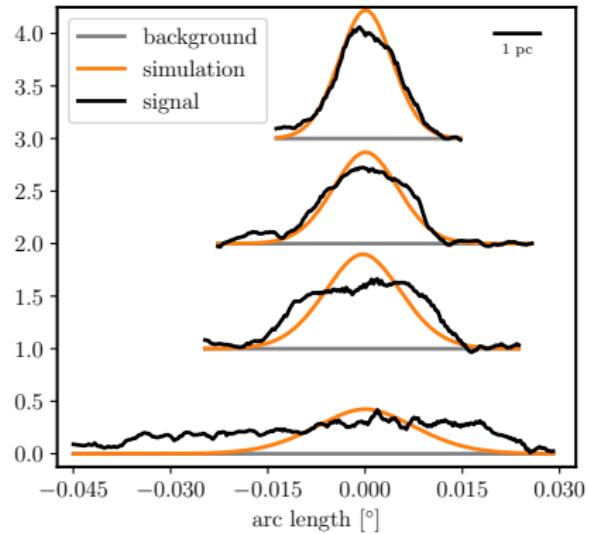


Radio synchrotron harps: testing CR propagation



Haywood+ (Nature, 2019)

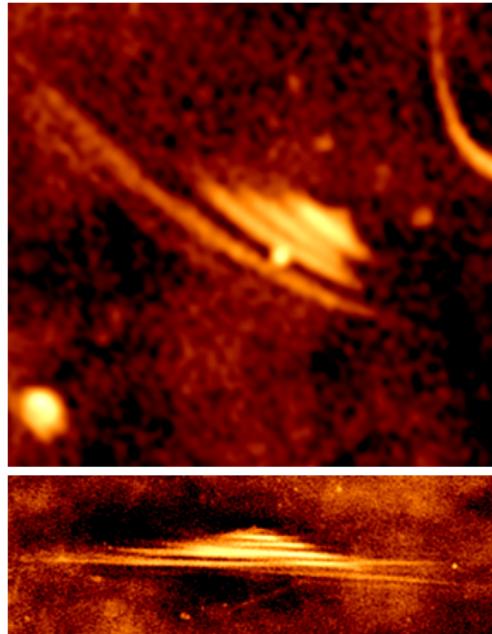
CR diffusion



Thomas, CP, Enßlin (2020)

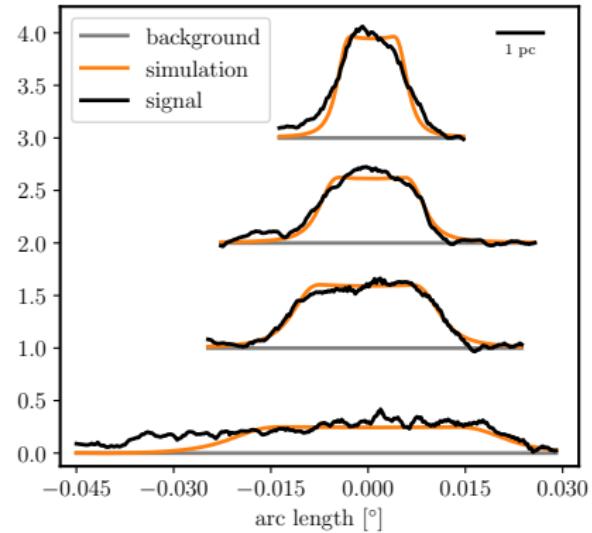


Radio synchrotron harps: testing CR propagation



Haywood+ (Nature, 2019)

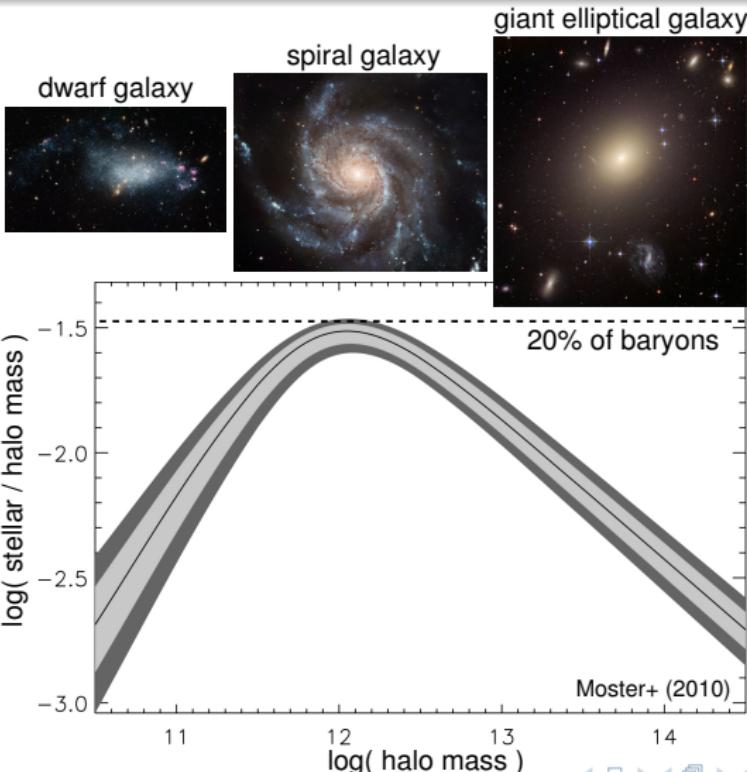
CR streaming and diffusion



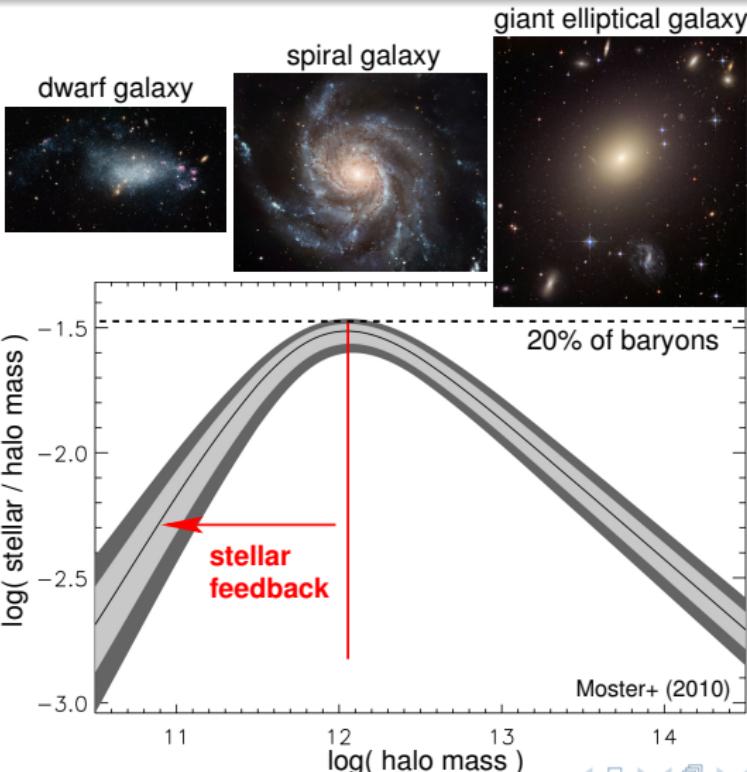
Thomas, CP, Enßlin (2020)



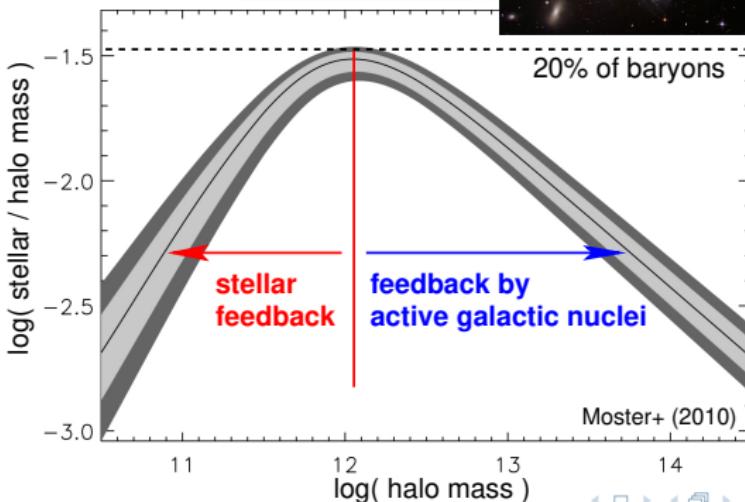
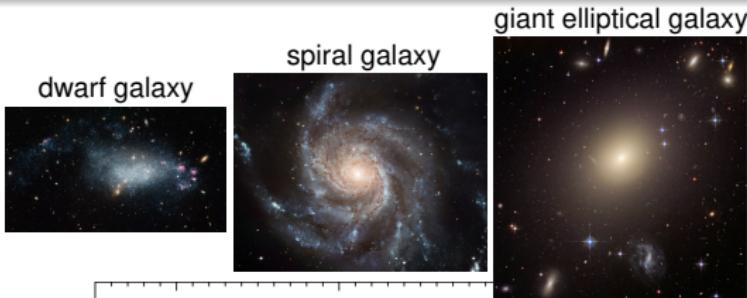
Puzzles in galaxy formation



Puzzles in galaxy formation



Puzzles in galaxy formation



AIP

How are galactic winds driven?



super wind in M82

- thermal pressure provided by supernovae or AGNs?
- radiation pressure and photoionization by massive stars and QSOs?
- cosmic-ray (CR) pressure and Alfvén wave heating of CRs accelerated at supernova shocks?



How are galactic winds driven?



NASA/JPL-Caltech/STScI/CXC/UofA

super wind in M82

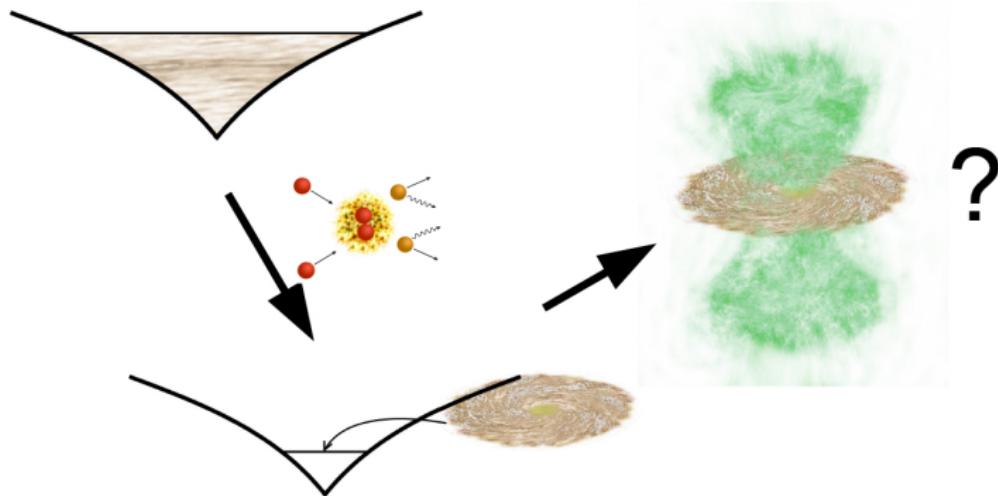
- thermal pressure provided by supernovae or AGNs?
- radiation pressure and photoionization by massive stars and QSOs?
- cosmic-ray (CR) pressure and Alfvén wave heating of CRs accelerated at supernova shocks?

observed energy equipartition between cosmic rays, thermal gas and magnetic fields

→ suggests self-regulated feedback loop with CR driven winds



1. Cosmic ray feedback in galaxy formation

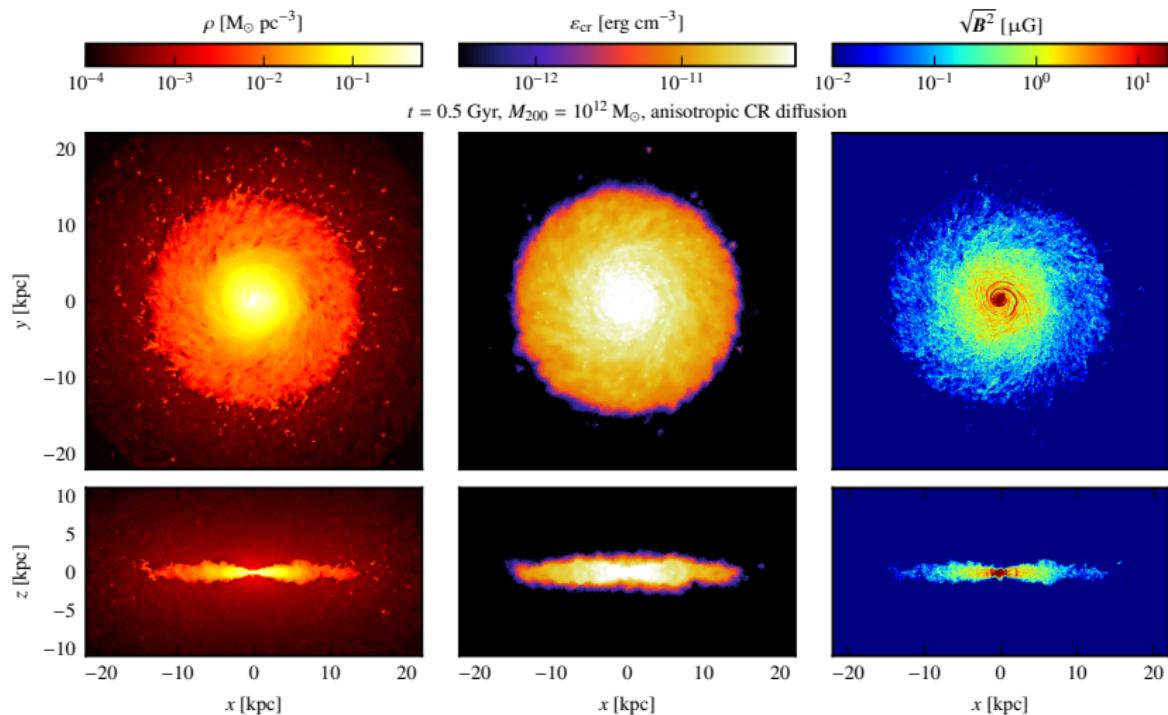


Pakmor, CP+ (2016), CP+ (2017b)

Galactic winds driven by CR diffusion in isolated disk galaxies

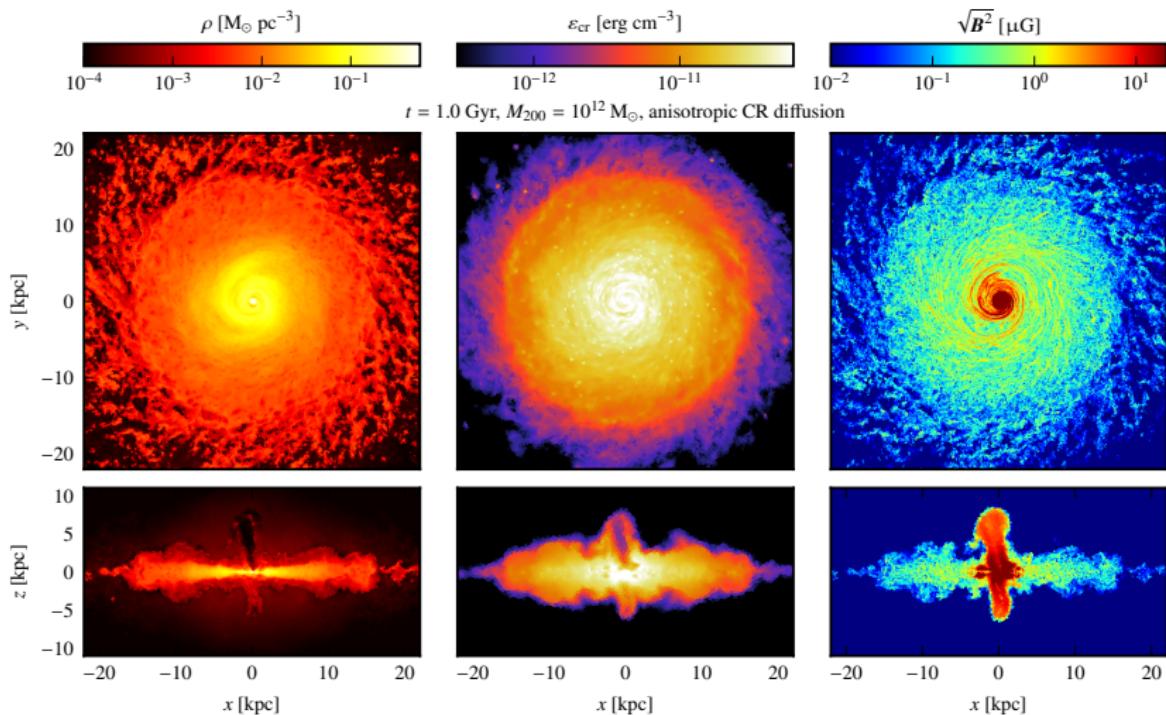
MHD + CR advection + anisotropic diffusion, $\{10^{10}, 10^{11}, 10^{12}\} M_{\odot}$

Simulation of Milky Way-like galaxy, $t = 0.5$ Gyr



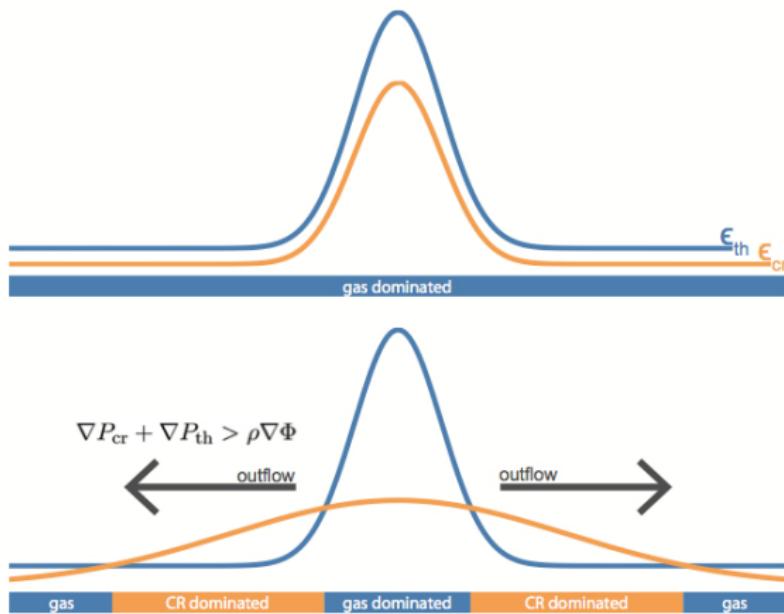
CP+ (2017b)

Simulation of Milky Way-like galaxy, $t = 1.0$ Gyr



CP+ (2017b)

Cosmic ray driven wind: mechanism

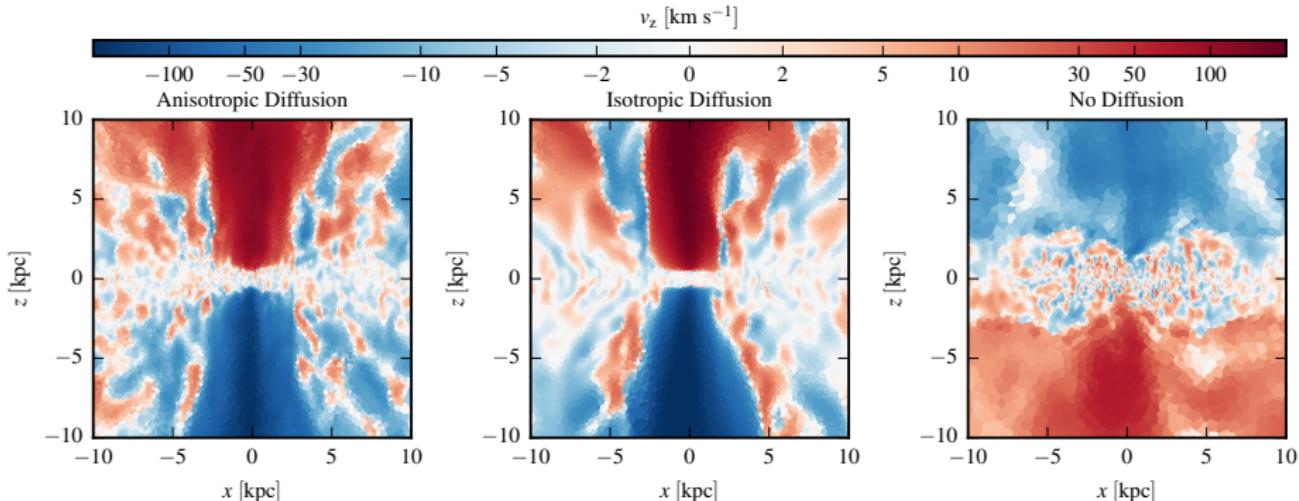


CR streaming in 3D simulations: Uhlig, CP+ (2012), Ruszkowski+ (2017)

CR diffusion in 3D simulations: Jubelgas+ (2008), Booth+ (2013), Hanasz+ (2013),
Salem & Bryan (2014), Pakmor, CP+ (2016), Simpson+ (2016), Girichidis+ (2016),
Dubois+ (2016), CP+ (2017b), Jacob+ (2018), ...



CR diffusion vs. advection

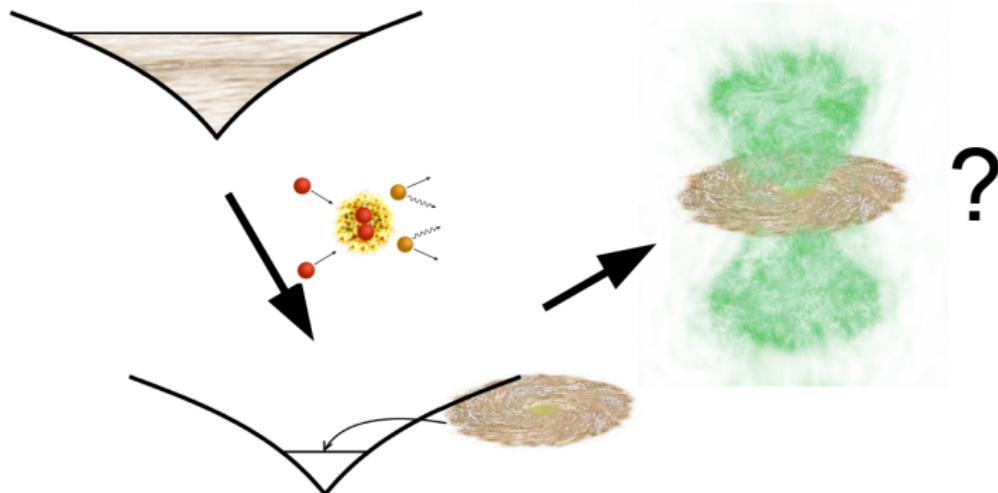


Pakmor, CP, Simpson, Springel (2016)

- CR diffusion launches powerful winds
- simulation without CR diffusion exhibits only weak fountain flows



2. Cosmic rays and non-thermal emission



Werhahn, CP+ (2021a, b, c), CP, Werhahn+ (2021)
Cosmic rays and non-thermal emission in simulated galaxies

MHD + CR advection + anisotropic diffusion: $\{10^{10}, 10^{11}, 10^{12}\} M_{\odot}$
steady-state spectra of CR protons, primary & secondary electrons

Steady-state cosmic ray spectra

- solve the steady-state equation in every cell for each CR population:

$$\frac{f(E)}{\tau_{\text{esc}}} - \frac{d}{dE} [f(E)b(E)] = q(E)$$

- protons: Coulomb, hadronic and escape losses (re-normalized to ε_{cr})
- electrons: Coulomb, bremsstr., IC, synchrotron and escape losses
 - primaries (re-normalized using $K_{\text{ep}} = 0.02$)
 - secondaries

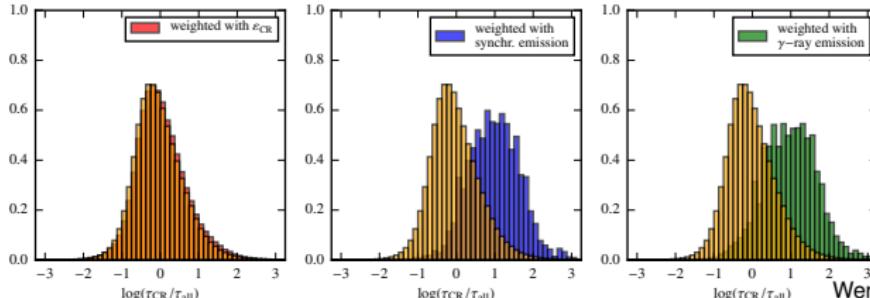


Steady-state cosmic ray spectra

- solve the steady-state equation in every cell for each CR population:

$$\frac{f(E)}{\tau_{\text{esc}}} - \frac{d}{dE} [f(E)b(E)] = q(E)$$

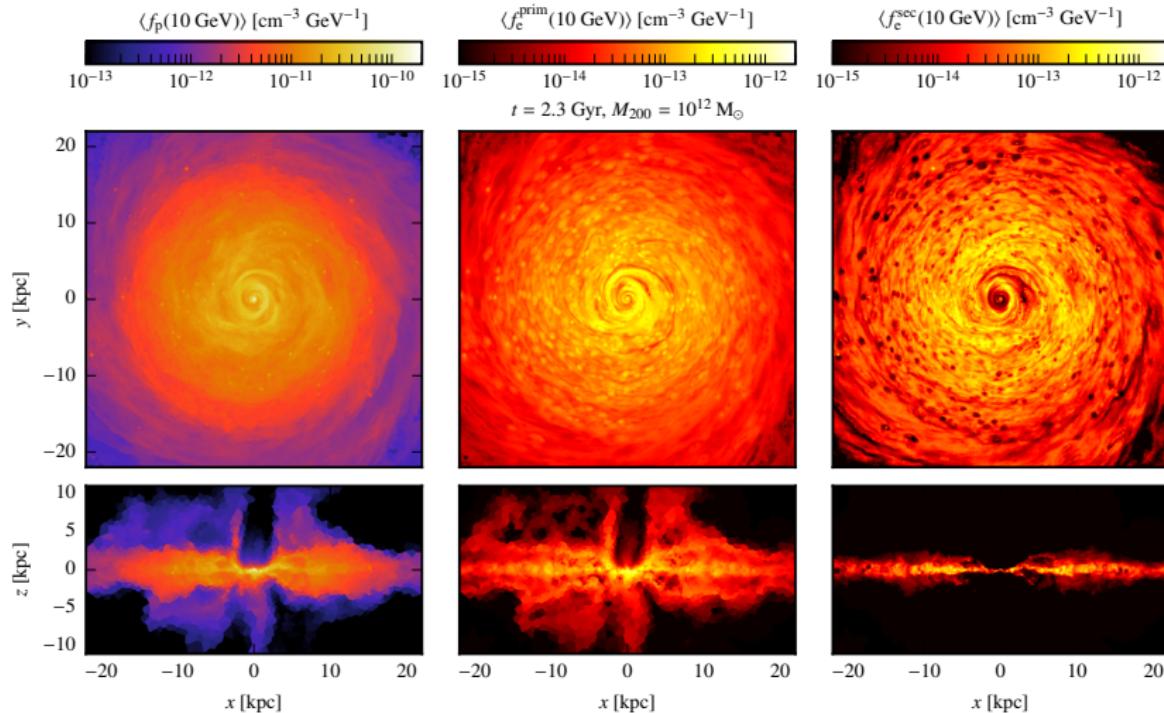
- protons:** Coulomb, hadronic and escape losses (re-normalized to ε_{cr})
- electrons:** Coulomb, bremsstr., IC, synchrotron and escape losses
 - primaries (re-normalized using $K_{\text{ep}} = 0.02$)
 - secondaries
- steady state assumption is fulfilled in disk** and in regions dominating the non-thermal emission but not at low densities, at SNRs and in outflows



Werhahn+ (2021a)

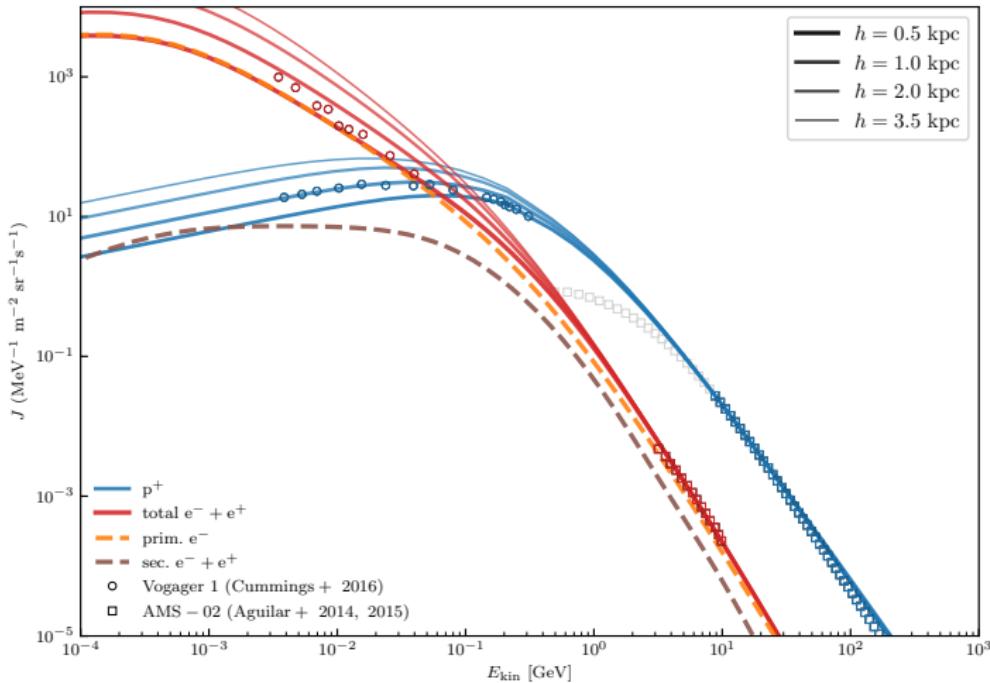


From a starburst galaxy to a Milky Way analogy



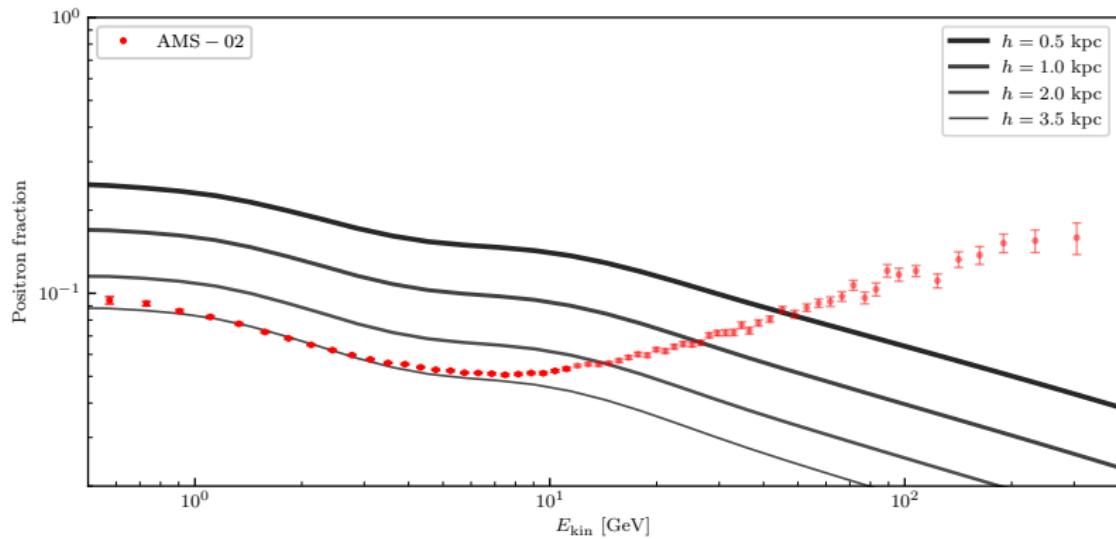
Werhahn, CP+ (2021b)

Comparing CR spectra to Voyager and AMS-02 data



Werhahn, CP+ (2021a)

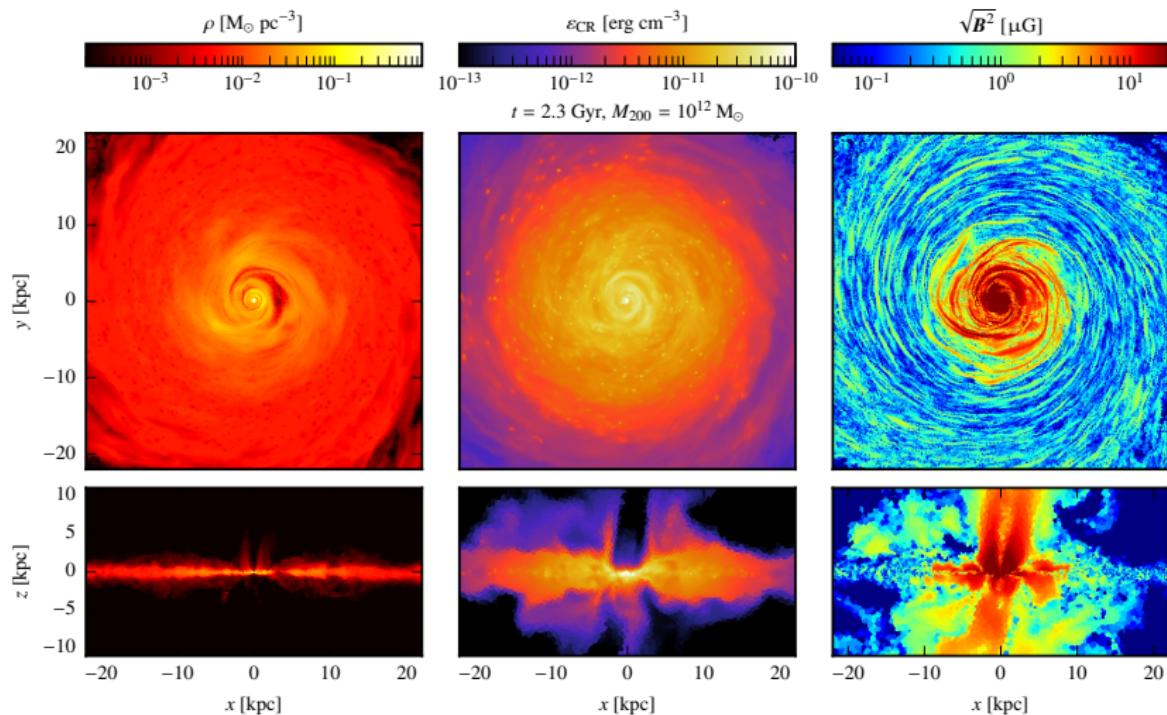
Comparing the positron fraction to AMS-02 data



Werhahn, CP+ (2021a)

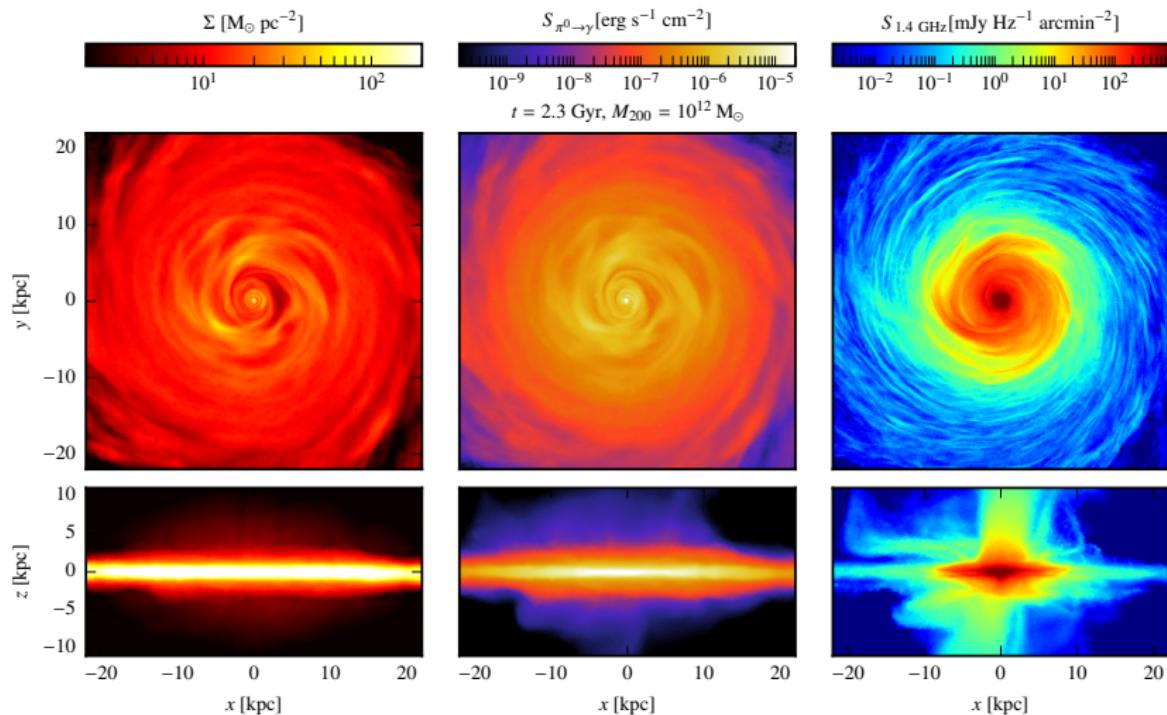


Simulation of a starburst galaxy



Werhahn, CP+ (2021b, c)

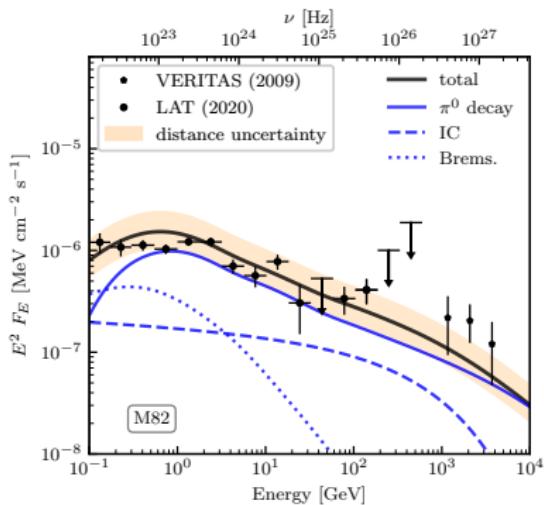
Simulation of a starburst galaxy



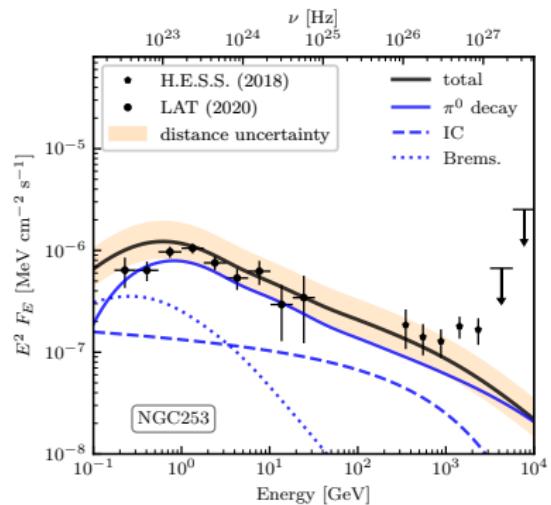
Werhahn, CP+ (2021b, c)

Gamma-ray spectra of starburst galaxies

Messier 82



NGC 253



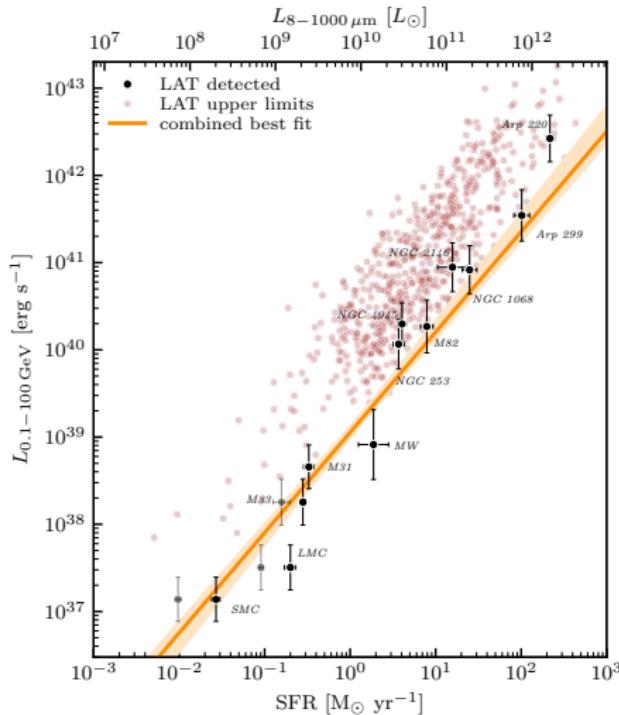
Werhahn, CP+ (2021b)

- gamma-ray spectra in starbursts **dominated by pion decay**
- CR protons propagate in **Kolmogorov turbulence**: $\kappa \propto E^{0.3}$



Far infra-red – gamma-ray correlation

Universal conversion: star formation → cosmic rays → gamma rays

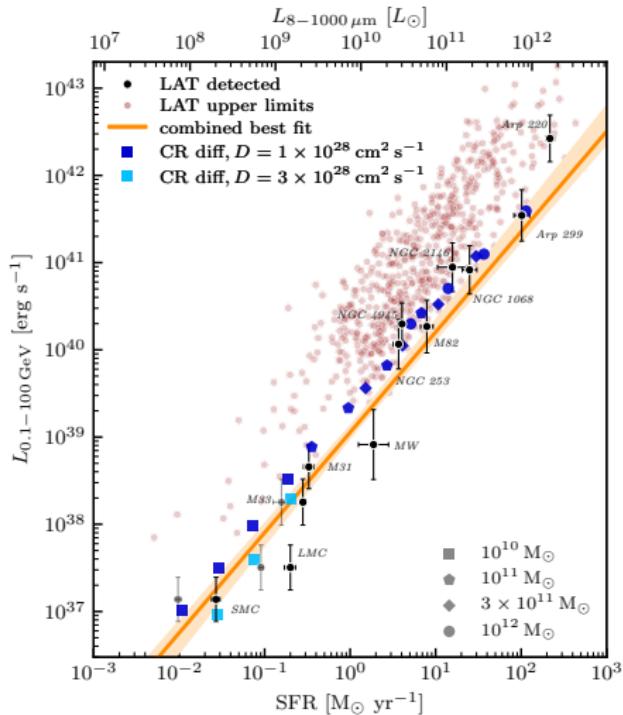


Ajello+ (2020)



Far infra-red – gamma-ray correlation

Universal conversion: star formation → cosmic rays → gamma rays

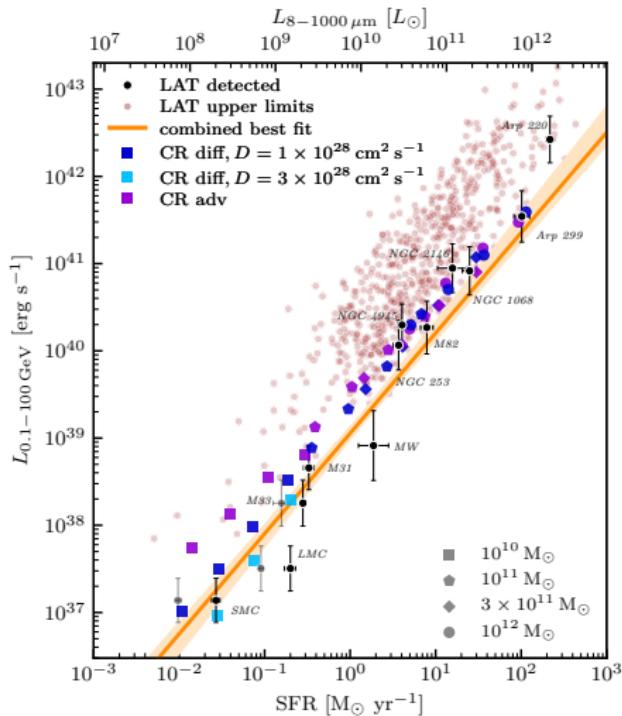


Werhahn, CP+ (2021b)



Far infra-red – gamma-ray correlation

Universal conversion: star formation → cosmic rays → gamma rays

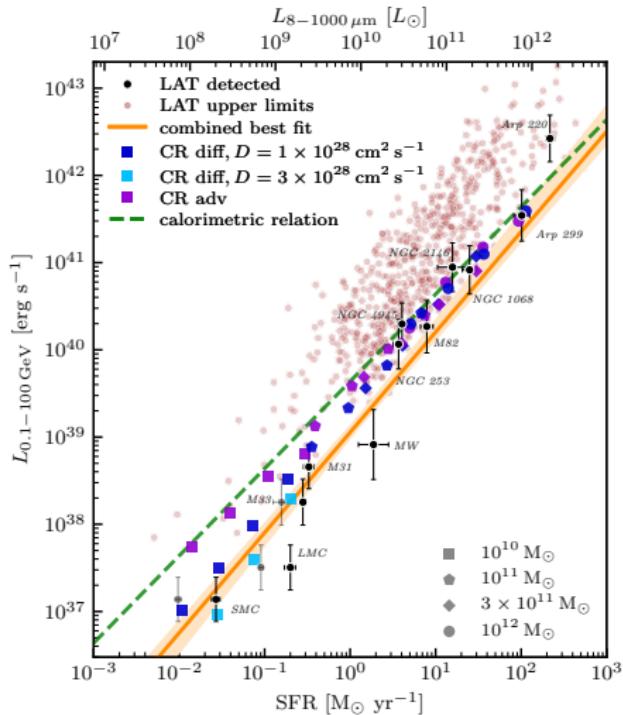


Werhahn, CP+ (2021b)



Far infra-red – gamma-ray correlation

Universal conversion: star formation → cosmic rays → gamma rays

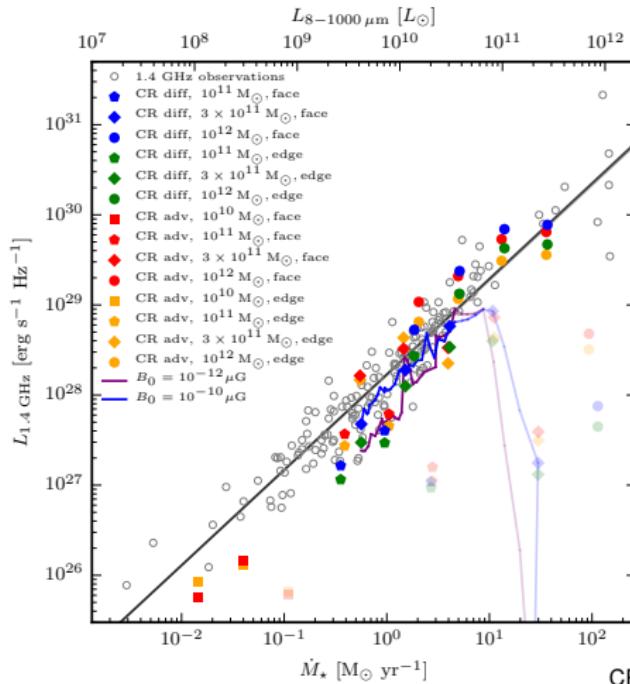


Werhahn, CP+ (2021b)



Far infra-red – radio correlation

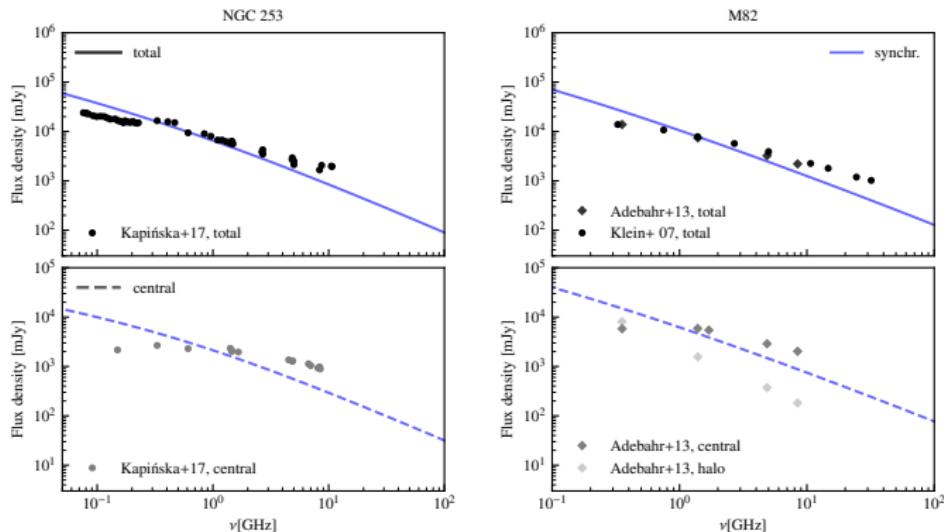
Universal conversion: star formation \rightarrow cosmic rays \rightarrow radio



CP, Werhahn+ (2021)



Radio spectra of starburst galaxies

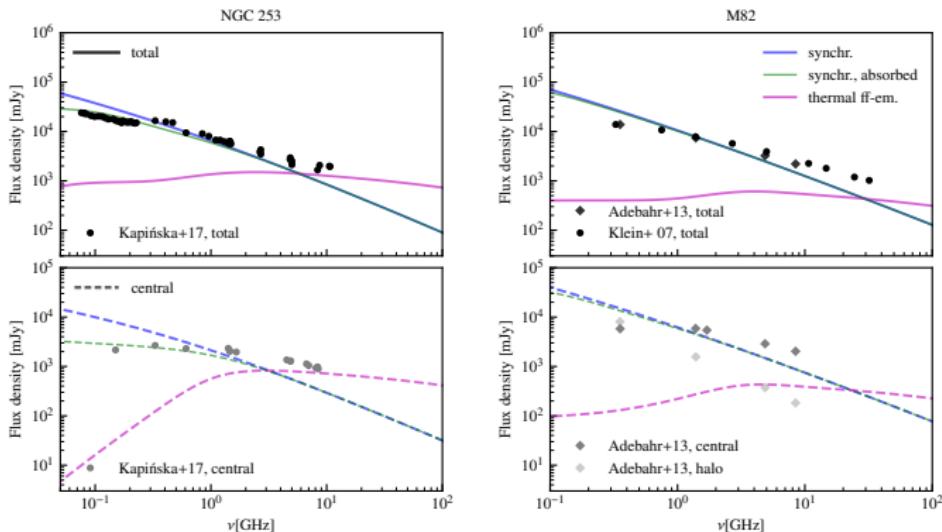


Werhahn, CP+ (2021c)

- steep synchrotron spectra demonstrate electron calorimetry



Radio spectra of starburst galaxies

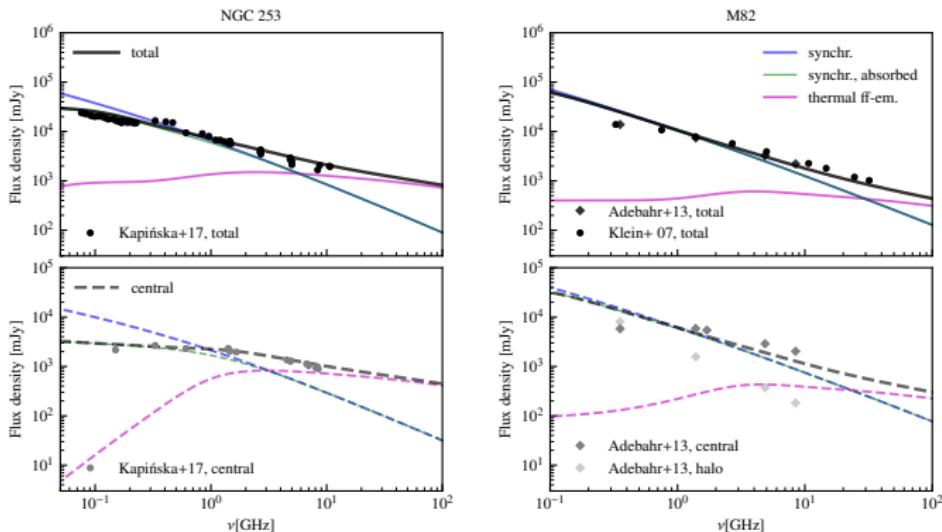


Werhahn, CP+ (2021c)

- steep synchrotron spectra demonstrate electron calorimetry
- free-free absorption/emission flatten spectra at low/high frequencies



Radio spectra of starburst galaxies



Werhahn, CP+ (2021c)

- steep synchrotron spectra demonstrate electron calorimetry
- free-free absorption/emission flatten spectra at low/high frequencies to match observations



Conclusions

CR-driven plasma instabilities:

- discovery of new *intermediate-scale instability*, which grows faster than the gyro-resonant instability
- implications for CR transport in galaxies, electron heating and injection into diffusive shock acceleration, and CR escape from acceleration sites



Conclusions

CR-driven plasma instabilities:

- discovery of new *intermediate-scale instability*, which grows faster than the gyro-resonant instability
- implications for CR transport in galaxies, electron heating and injection into diffusive shock acceleration, and CR escape from acceleration sites

CR transport in galaxies:

- novel theory of CR transport mediated by Alfvén waves and coupled to magneto-hydrodynamics
- synchrotron harps: CR streaming dominates over diffusion



Conclusions

CR-driven plasma instabilities:

- discovery of new *intermediate-scale instability*, which grows faster than the gyro-resonant instability
- implications for CR transport in galaxies, electron heating and injection into diffusive shock acceleration, and CR escape from acceleration sites

CR transport in galaxies:

- novel theory of CR transport mediated by Alfvén waves and coupled to magneto-hydrodynamics
- synchrotron harps: CR streaming dominates over diffusion

CR-induced signatures in galaxies

- CR feedback drives galactic winds & slows down star formation
- 3D MHD simulations reproduce CR, gamma-ray and radio spectra & FIR radio/gamma-ray correlations



AIP

CRAGSMAN: The Impact of Cosmic RAys on Galaxy and CluSter ForMAtion



This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (grant agreement No CRAGSMAN-646955).



Literature for the talk – 1

Cosmic ray driven instabilities:

- Shalaby, Thomas, Pfrommer, *A new cosmic ray-driven instability*, 2021, ApJ.

Cosmic ray transport:

- Pfrommer, Pakmor, Schaal, Simpson, Springel, *Simulating cosmic ray physics on a moving mesh*, 2017a, MNRAS.
- Thomas & Pfrommer, *Cosmic-ray hydrodynamics: Alfvén-wave regulated transport of cosmic rays*, 2019, MNRAS.
- Thomas, Pfrommer, Enßlin, *Probing cosmic ray transport with radio synchrotron harps in the Galactic center*, 2020, ApJL.
- Thomas, Pfrommer, Pakmor, *A finite volume method for two-moment cosmic-ray hydrodynamics on a moving mesh*, 2021, MNRAS
- Thomas & Pfrommer, *Comparing different closure relations for cosmic ray hydrodynamics*, 2021, MNRAS



Literature for the talk – 2

Cosmic ray feedback in galaxies:

- Pakmor, Pfrommer, Simpson, Springel, *Galactic winds driven by isotropic and anisotropic cosmic ray diffusion in isolated disk galaxies*, 2016, ApJL.

Cosmic rays and non-thermal emission in galaxies:

- Pfrommer, Pakmor, Simpson, Springel, *Simulating gamma-ray emission in star-forming galaxies*, 2017b, ApJL.
- Werhahn, Pfrommer, Girichidis, Puchwein, Pakmor, *Cosmic rays and non-thermal emission in simulated galaxies. I. Electron and proton spectra explain Voyager-1 data*, 2021a, MNRAS
- Werhahn, Pfrommer, Girichidis, Winner, *Cosmic rays and non-thermal emission in simulated galaxies. II. γ -ray maps, spectra and the far infrared- γ -ray relation*, 2021b, MNRAS
- Werhahn, Pfrommer, Girichidis, *Cosmic rays and non-thermal emission in simulated galaxies. III. probing cosmic ray calorimetry with radio spectra and the FIR-radio correlation*, 2021c
- Pfrommer, Werhahn, Pakmor, Girichidis, Simpson, *Simulating radio synchrotron emission in star-forming galaxies: small-scale magnetic dynamo and the origin of the far infrared-radio correlation*, 2021

