

Cosmic-ray driven plasma instabilities

Christoph Pfrommer

Collaborators: Mohamad Shalaby, Rouven Lemmerz

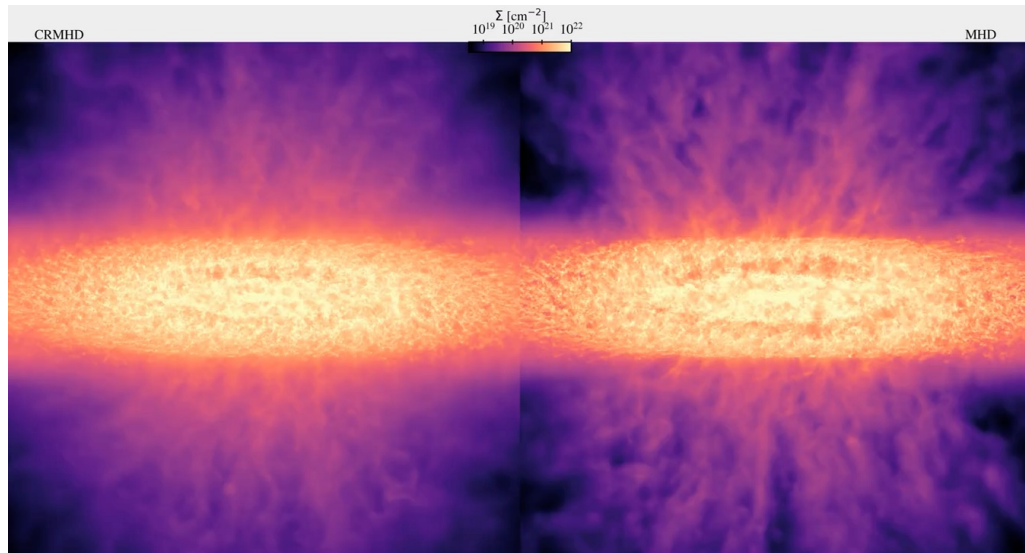
AIP Potsdam

2nd July 2025

Cosmic Rays Impact Galaxies

with CRs

without CRs



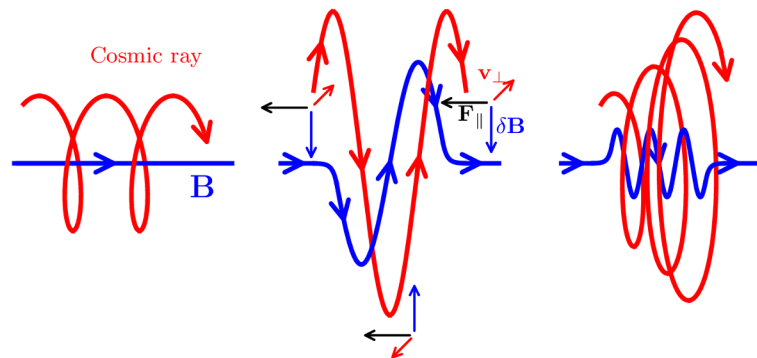
Cosmic rays

- drive galactic winds
- regulate star formation
- amplify magnetic fields through microphysical interactions

Thomas, CP, Pakmor 2025

Connecting the Scales

Gyroresonant instabilities



Jacob & CP

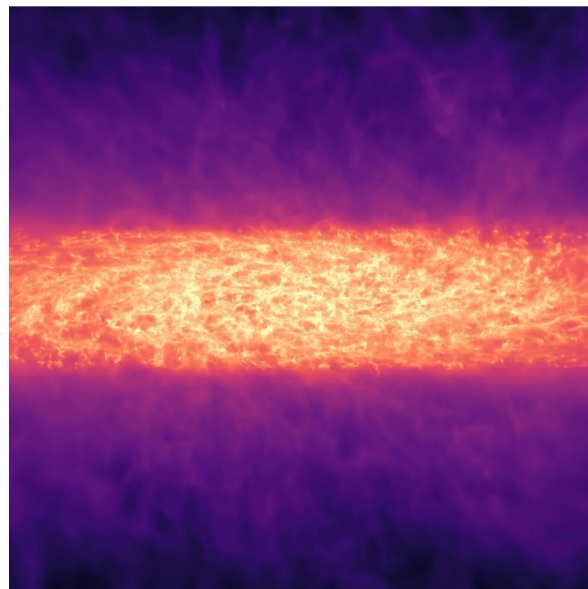
Scale \sim AU

Impact of cosmic ray interactions

CR model



Diffusion
Mean velocity
Pressure



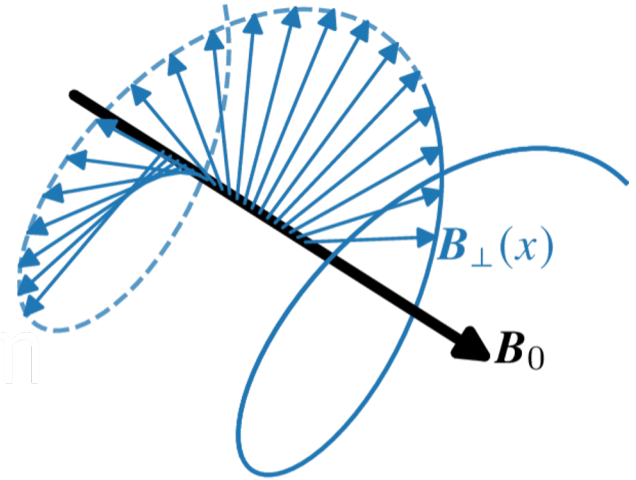
Scale \gtrsim kpc (10^8 AU)

What is Gyroresonance?

Plane wave: $\exp(-ik(x - v_{\text{wave}}t))$



Leibniz-Institut für
Astrophysik Potsdam



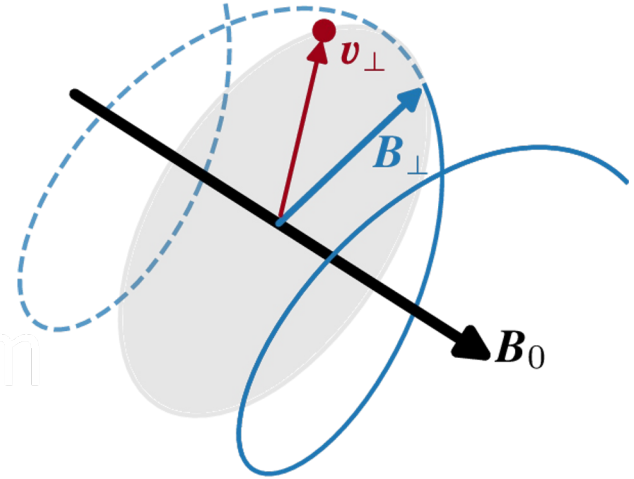
Lemmerz+ 2025

What is Gyroresonance?

Plane wave: $\exp(-ik(x - v_{\text{wave}}t))$

Cosmic ray: v_{\parallel} movement along B_0

Ω_{cr} gyration frequency



Lemmerz+ 2025

What is Gyroresonance?

Plane wave: $\exp(-ik(x - v_{\text{wave}}t))$

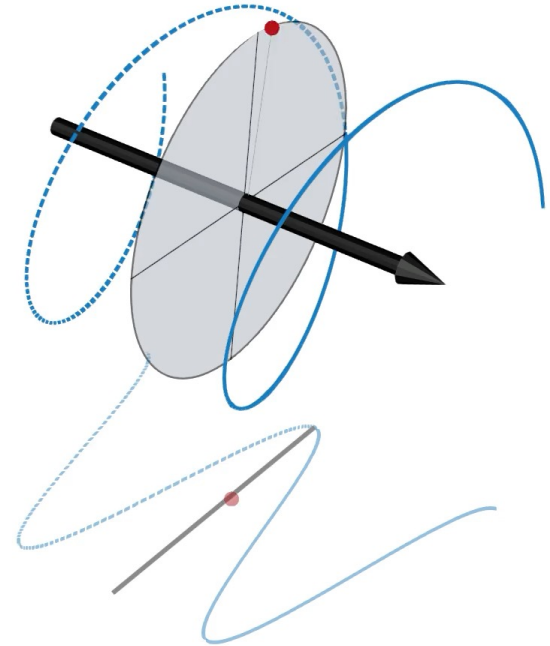
Cosmic ray: v_{\parallel} movement along B_0
 Ω_{cr} gyration frequency

Resonance condition:

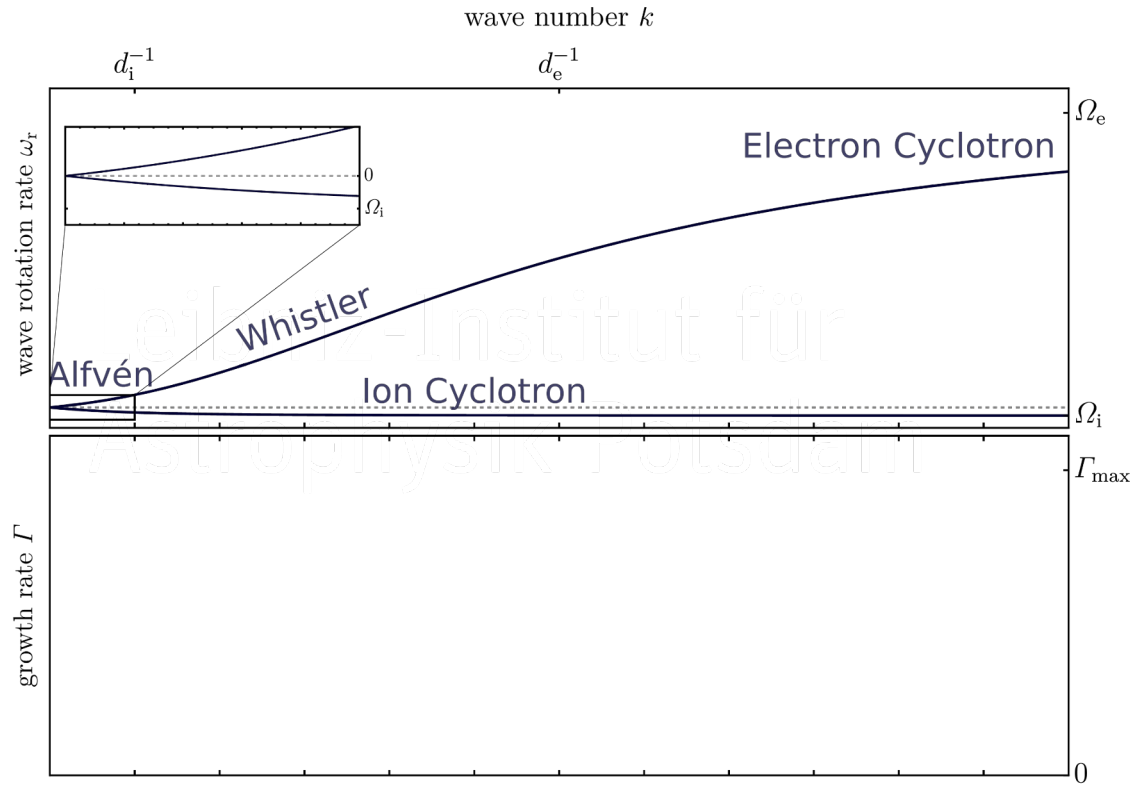
$$\underbrace{\Omega_{\text{cr}}}_{\text{Gyration}} + \underbrace{kv_{\parallel}}_{\text{Dopplershift}} = \underbrace{kv_{\text{wave}}}_{\text{wave frequency}}$$

Resonant wave appears **static** to CR

Wave frame

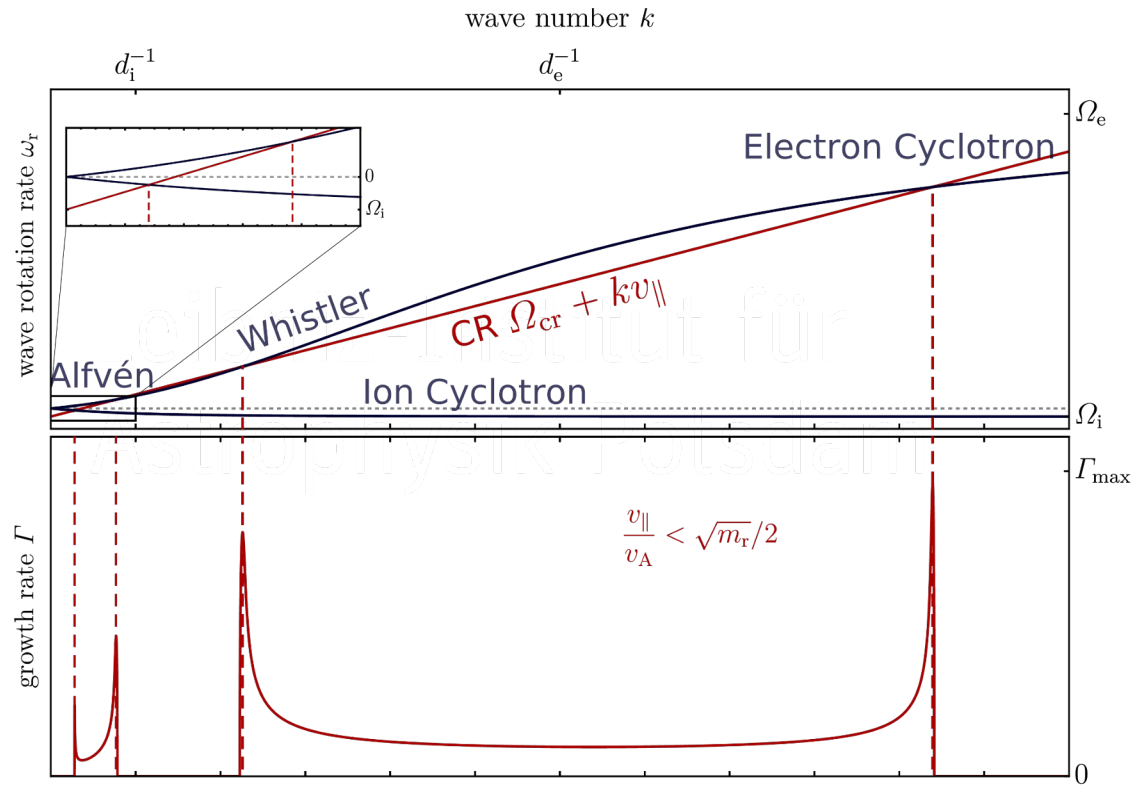


Gyroresonance in Dispersion Relation



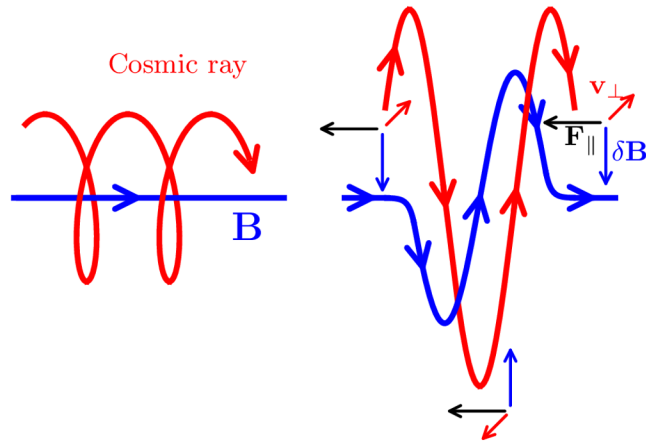
Shalaby+ 2023

Gyroresonance in Dispersion Relation



Shalaby+ 2023

Particle-Wave Interactions



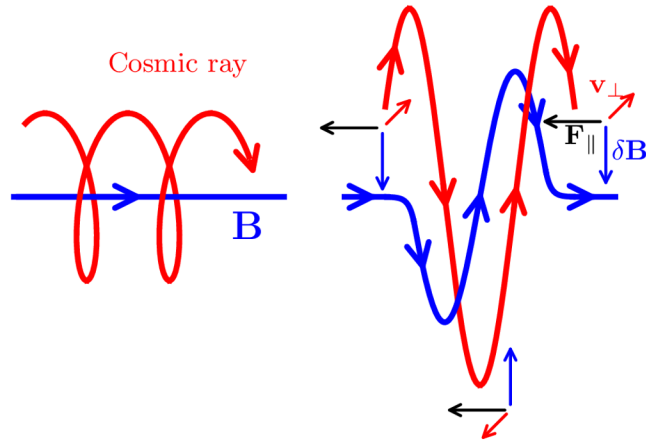
- electric fields vanish in the Alfvén wave frame:

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

für
sdam

Jacob & CP

Particle-Wave Interactions



- electric fields vanish in the Alfvén wave frame:

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

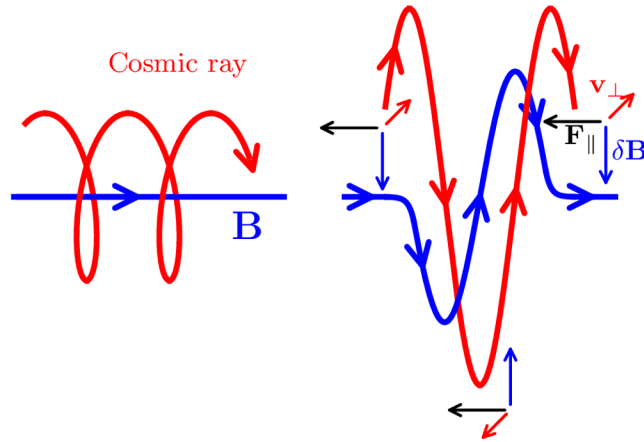
- thus energy is conserved:

$$p^2 = p_{\parallel}^2 + p_{\perp}^2 = \text{const.}$$

für
sdam

Jacob & CP

Particle-Wave Interactions



- electric fields vanish in the Alfvén wave frame:

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

- thus energy is conserved:

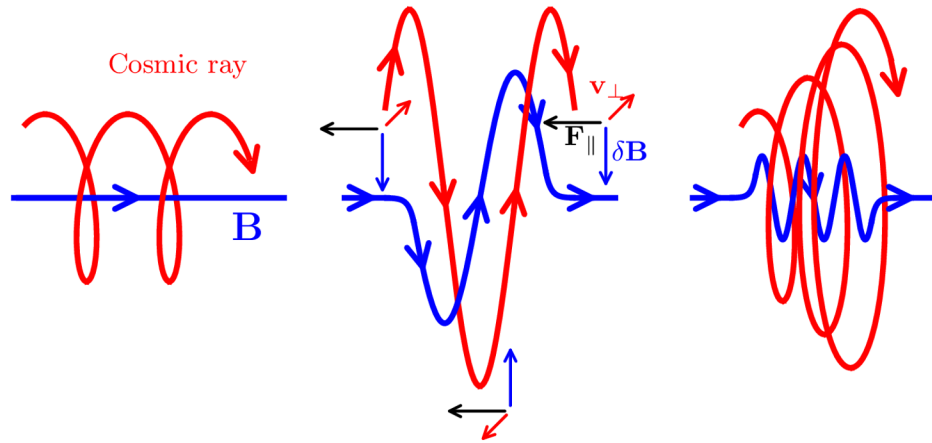
$$p^2 = p_{\parallel}^2 + p_{\perp}^2 = \text{const.}$$

- work out Lorentz force:

$$\mathbf{F}_L = q \frac{\mathbf{v} \times \mathbf{B}}{c}$$

Jacob & CP

Particle-Wave Interactions



Jacob & CP

- electric fields vanish in the Alfvén wave frame:

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

- thus energy is conserved:

$$p^2 = p_{\parallel}^2 + p_{\perp}^2 = \text{const.}$$

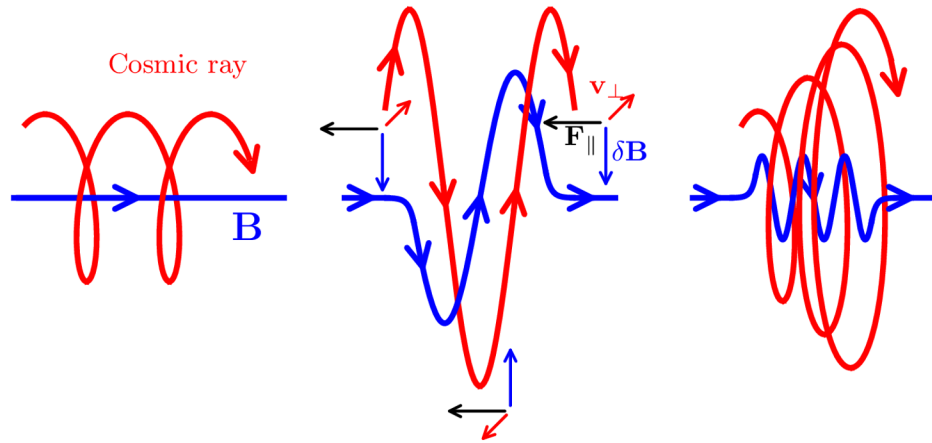
- work out Lorentz force:

$$\mathbf{F}_L = q \frac{\mathbf{v} \times \mathbf{B}}{c}$$

- which changes p_{\parallel} and thereby the pitch-angle cosine

$$\mu = \cos \theta = \frac{\mathbf{B}}{|\mathbf{B}|} \cdot \frac{\mathbf{p}}{|\mathbf{p}|}$$

Particle-Wave Interactions



Jacob & CP

=> But why do waves grow?

- electric fields vanish in the Alfvén wave frame:

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

- thus energy is conserved:

$$p^2 = p_{\parallel}^2 + p_{\perp}^2 = \text{const.}$$

- work out Lorentz force:

$$\mathbf{F}_L = q \frac{\mathbf{v} \times \mathbf{B}}{c}$$

- which changes p_{\parallel} and thereby the pitch-angle cosine

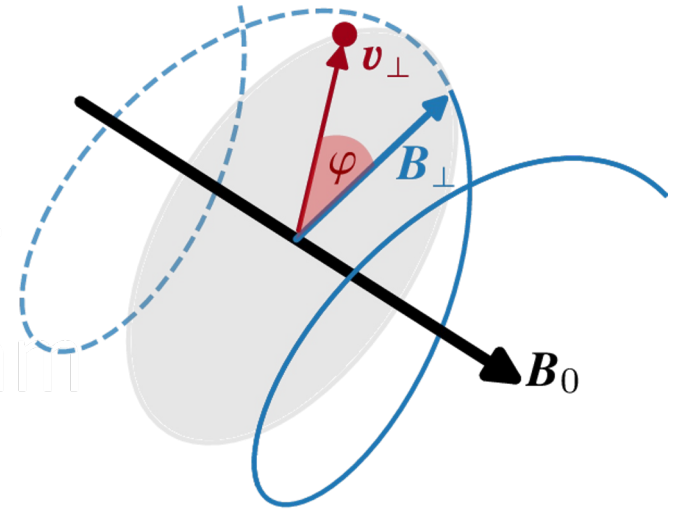
$$\mu = \cos \theta = \frac{\mathbf{B}}{|\mathbf{B}|} \cdot \frac{\mathbf{p}}{|\mathbf{p}|}$$

From Particle-Wave to Wave-Wave

- Goal: understand collective behaviour of many CRs



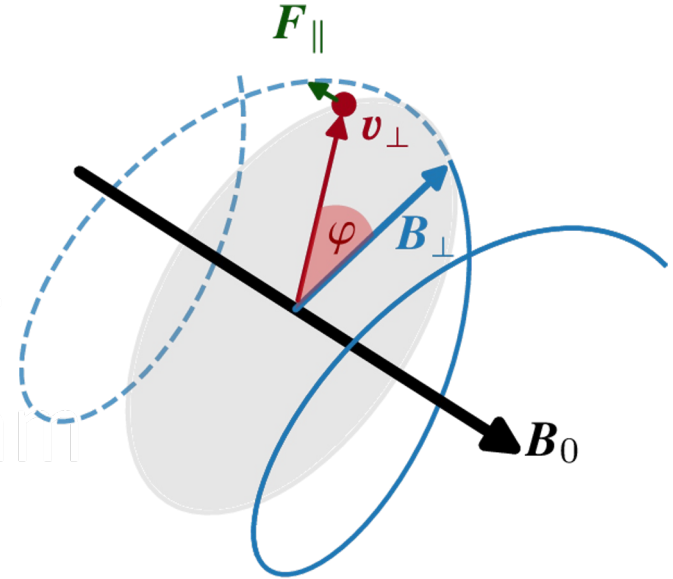
Leibniz-Institut für
Astrophysik Potsdam



Lemmerz+ 2025

From Particle-Wave to Wave-Wave

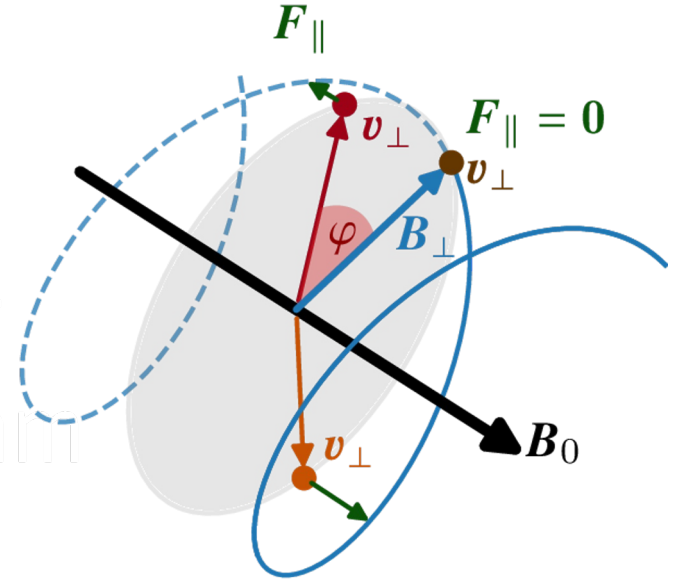
- Goal: understand collective behaviour of many CRs
- v_{\parallel} changes (Lorentz force)
 - acceleration depends on $\sim \sin(\varphi)$



Lemmerz+ 2025

From Particle-Wave to Wave-Wave

- Goal: understand collective behaviour of many CRs
- v_{\parallel} changes (Lorentz force)
 - acceleration depends on $\sim \sin(\varphi)$
- CRs align rotational phase with wave

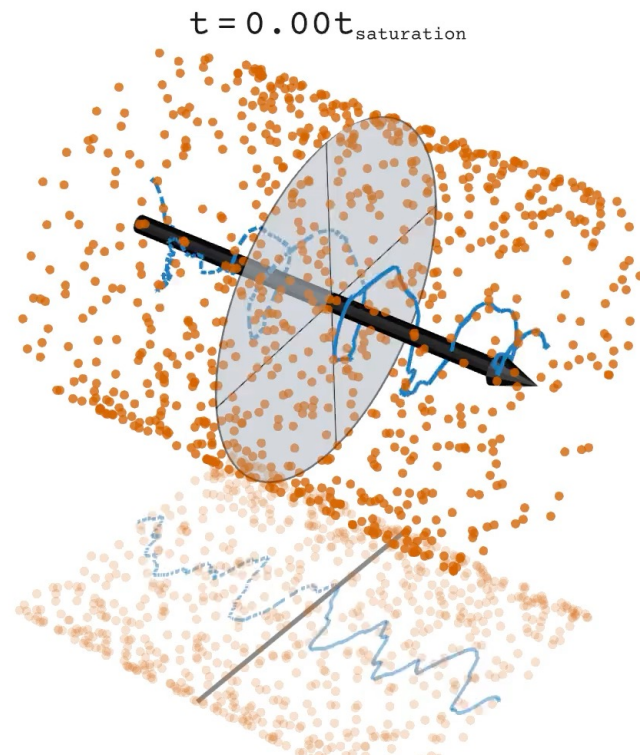


Lemmerz+ 2025

From Particle-Wave to Wave-Wave

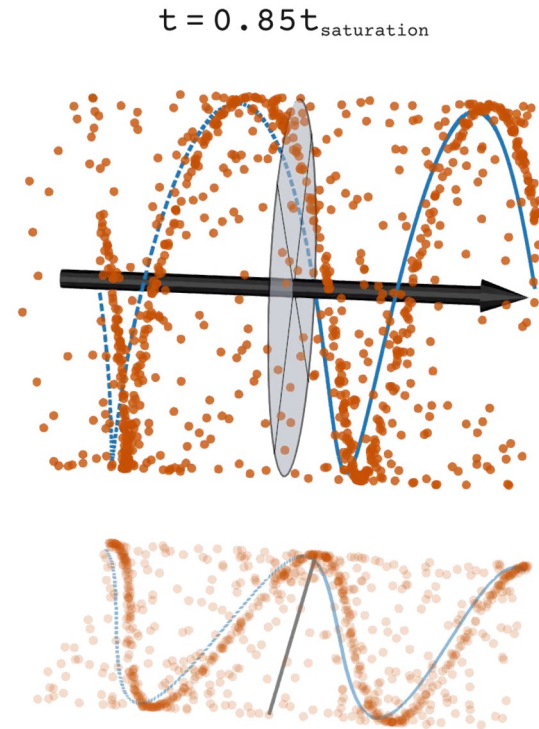
- Goal: understand collective behaviour of many CRs
- v_{\parallel} changes (Lorentz force)
 - acceleration depends on $\sim \sin(\varphi)$
- CRs align rotational phase with wave
- CRs **bunch** up

fluid-PIC simulation (Lemmerz+ 2025)



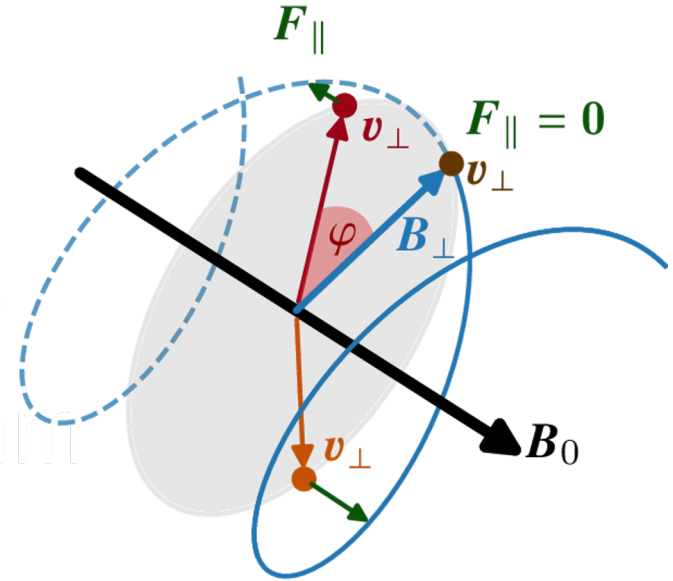
From Particle-Wave to Wave-Wave

- Goal: understand collective behaviour of many CRs
- v_{\parallel} changes (Lorentz force)
 - acceleration depends on $\sim \sin(\varphi)$
- CRs align rotational phase with wave
- CRs **bunch** up
- CR current wave interacts with EM wave



From Particle-Wave to Wave-Wave

- Goal: understand collective behaviour of many CRs
- v_{\parallel} changes (Lorentz force)
 - acceleration depends on $\sim \sin(\varphi)$
- CRs align rotational phase with wave
- CRs **bunch** up
- CR current wave interacts with EM wave

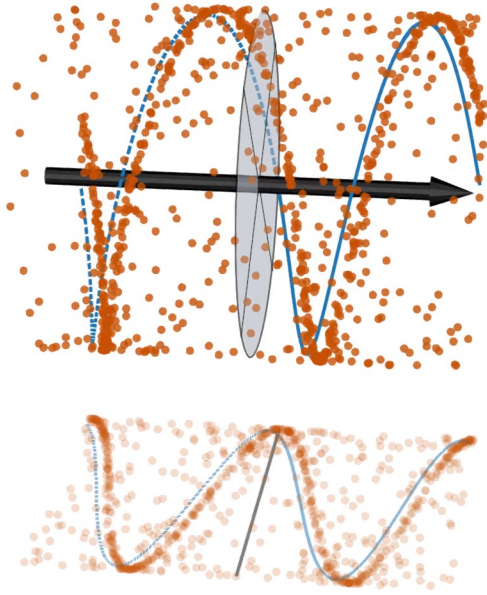


Lemmerz+ 2025

Gyroresonance with different waves

fluid-PIC simulation

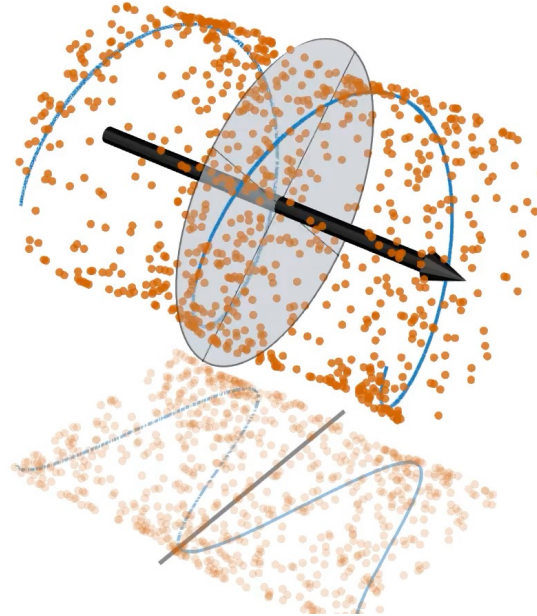
$t = 0.85 t_{\text{saturation}}$



Forward Alfvén, Whistler

Christoph Pfrommer

$t = 0.70 t_{\text{saturation}}$



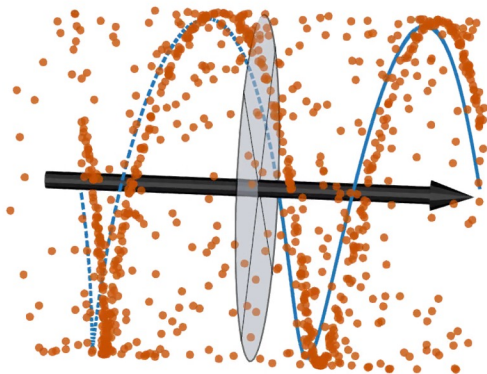
Backward Alfvén

CR Resonant Instabilities

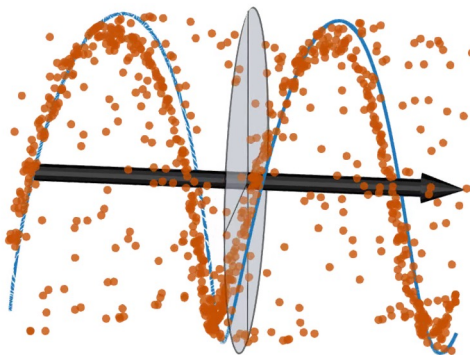
July 2, 2025

Gyroresonance with different waves

$t = 0.85 t_{\text{saturation}}$



$t = 0.85 t_{\text{saturation}}$



Bunching theory

- Bunching in gyrophase
- Biased scattering, favors wave growth

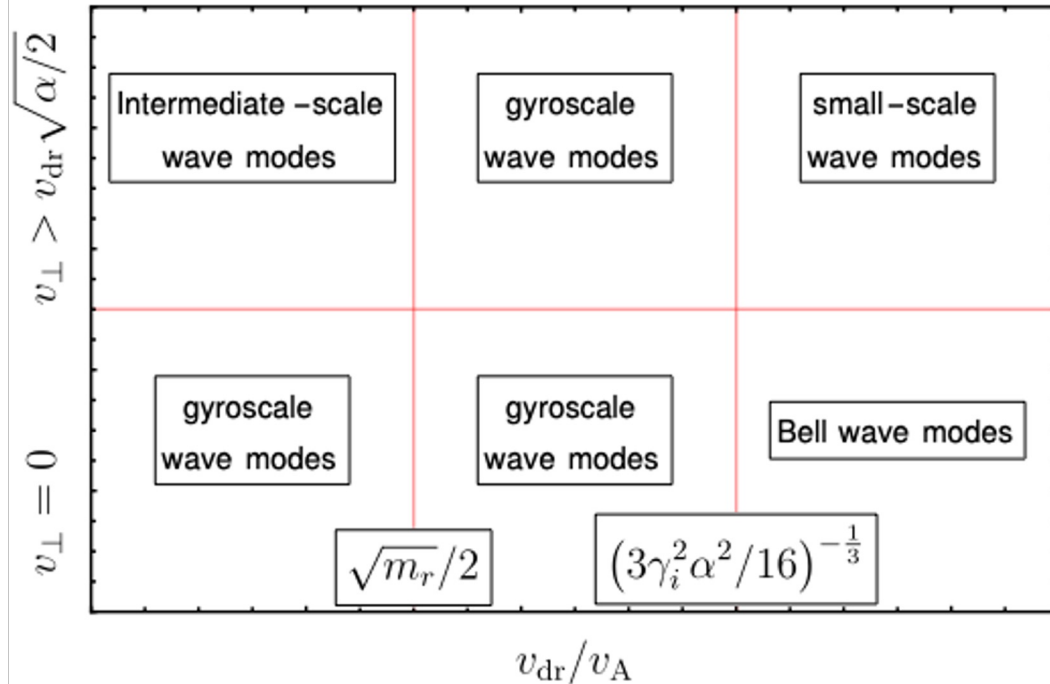
Traditional, Quasilinear theory

- Assumes uniform φ
- Diffusive scattering, no backward wave

Forward Alfvén, Whistler

Backward Alfvén

Graphical classification of CR instabilities

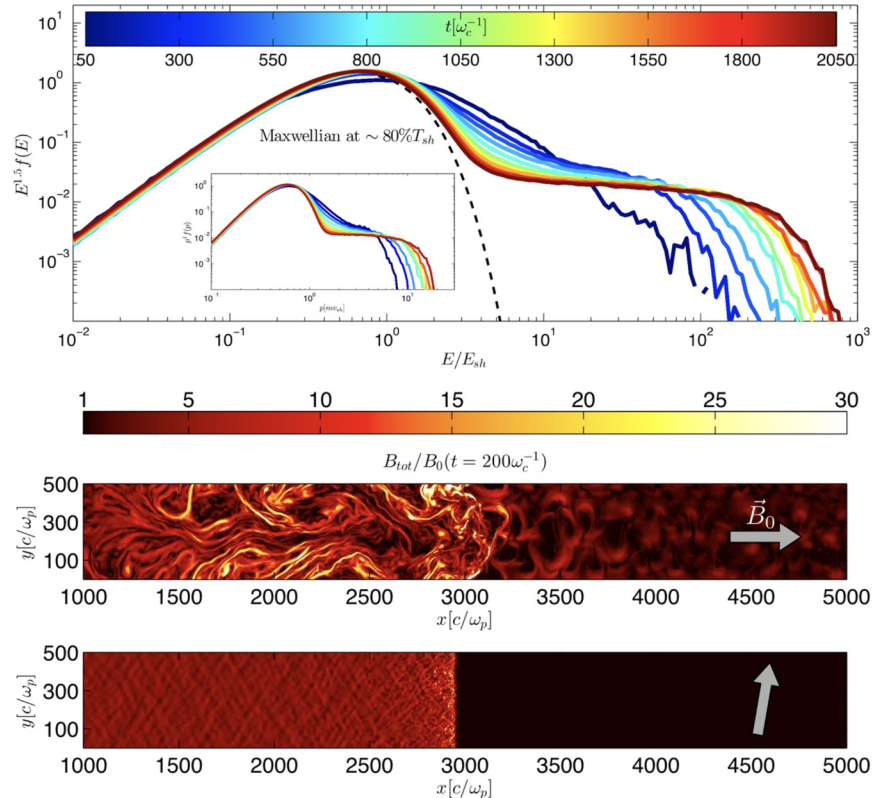


CR-driven instabilities in the linear regime (low density, gyrotropic CRs with a cold momentum distribution):

- the fastest wave modes depend on the CR flux and pitch angle
- because CRs will typically have a finite pitch angle, the typical dominant unstable wave modes occupy the top region

Shalaby+ 2021

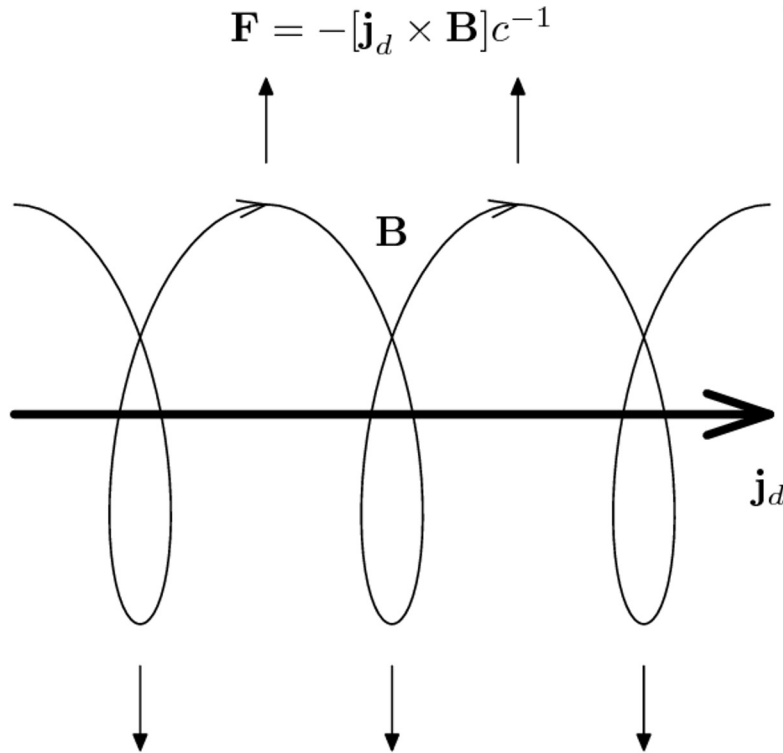
Non-resonant hybrid (Bell's) instability



Hybrid-PIC simulation of CR ion acceleration at a collisionless, non-relativistic strong shock (Caprioli & Spitkowski 2014).

- top panel: downstream ion energy spectrum of a quasi-parallel shock, color coded by different times – thermal Maxwellian & CR power law that shows an increasing maximum energy with time
- bottom panels: magnitude of the total magnetic field for strong shocks with different obliquity, implying that magnetic field amplification and CR acceleration only works in quasi-parallel shocks

Non-resonant hybrid (Bell's) instability



Visualization of the underlying principle of Bell's streaming instability:

- The CR current, \mathbf{j}_d , induces a return current in the background electrons, $-\mathbf{j}_d$, which amplifies a helical magnetic perturbation and stretches it via the Lorentz force $\mathbf{F} = -(\mathbf{j}_d \times \mathbf{B})c^{-1}$ (Zirakashvili+ 2008)
- Saturation once CRs get magnetized:

$$\varepsilon_{B,\text{sat}} \sim \frac{1}{2} \frac{v_s}{c} \varepsilon_{\text{cr}}$$

Main questions

→ Physics of instability saturation

- ◆ Saturation as growth = damping or through particle trapping in Lorentz force potential?
- ◆ Growth: which instability dominates? Forward/backward Alfvén, Whistler modes? Is MHD sufficient or do we need full dispersion relation?
- ◆ Damping: Ion-neutral, nonlinear Landau damping, turbulent damping?

→ Effect of inhomogeneities in ρ , \mathbf{B} , ...

→ Multi-D effects: turbulence, cascades, oblique waves